



S J P N Trust's

**Hirasugar Institute of Technology, Nidasoshi.**

*Inculcating Values, Promoting Prosperity*

Approved by AICTE, Recognized by Govt. of Karnataka and Affiliated to VTU Belagavi

ECE Dept.

EE

III Sem

2017-18

**Department of Electronics & Communication Engg.**

**Course : Engineering Electromagnetics -15EC36.**

**Sem.: 3<sup>rd</sup> (2017-18)**

**Course Coordinator:**

**Prof. S. S. KAMATE**

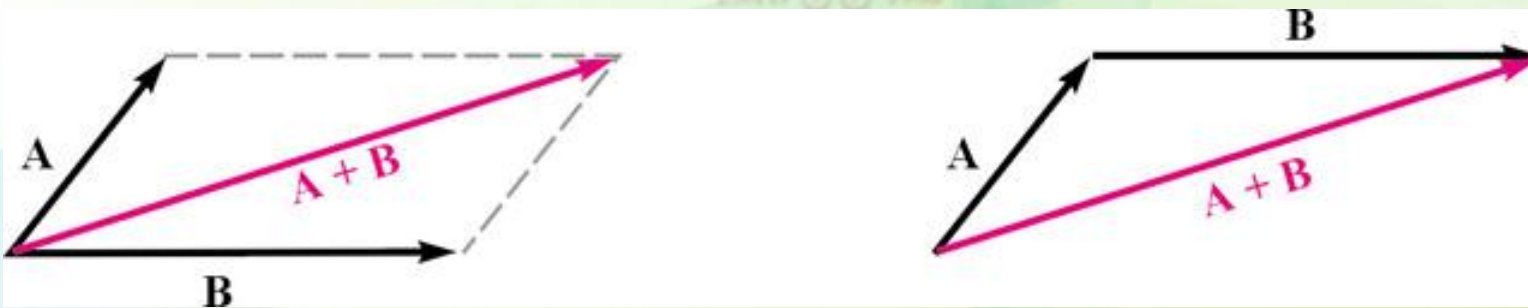
# Vector Analysis

## Scalars and Vectors

Scalar Fields (temperature)

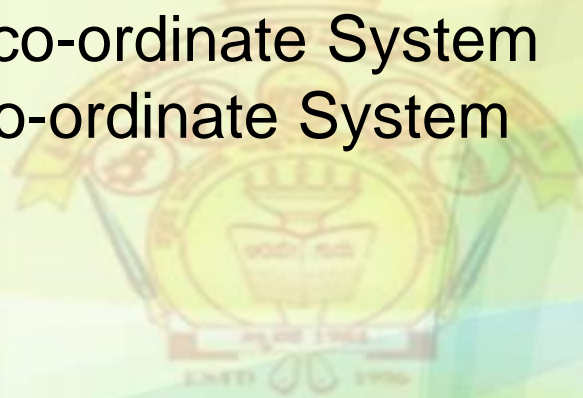
Vector Fields (gravitational, magnetic)

## Vector Algebra

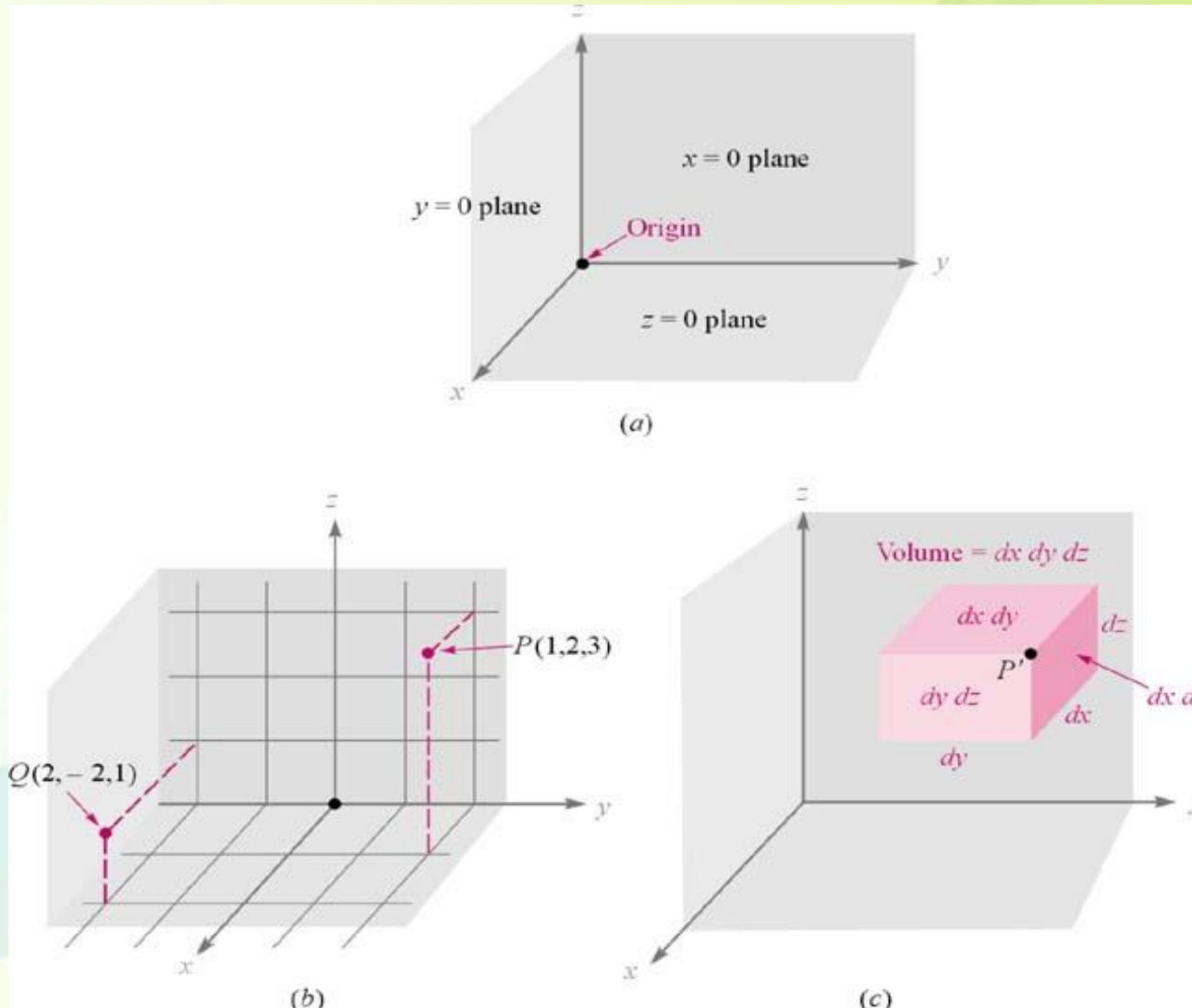


# There are three co-ordinate systems

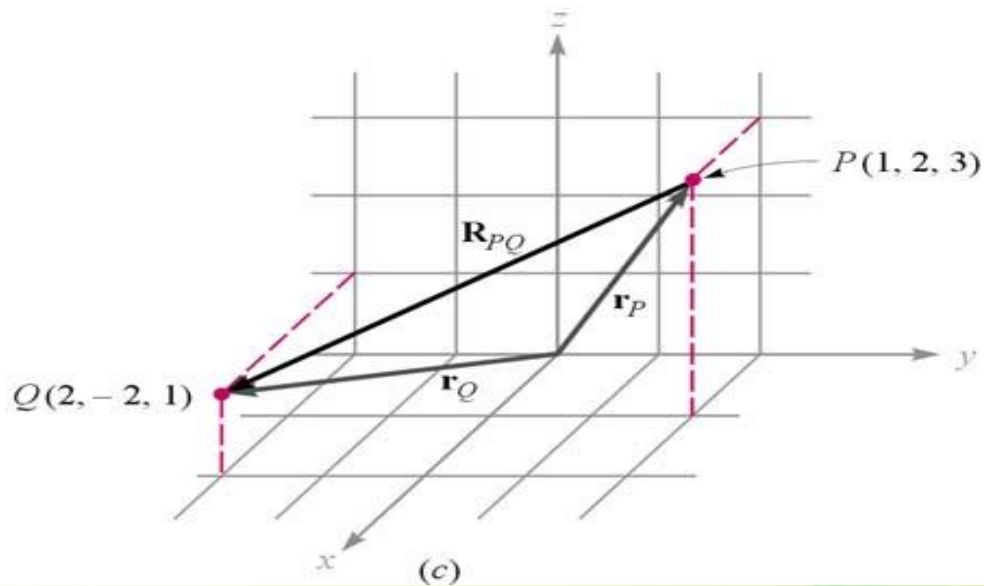
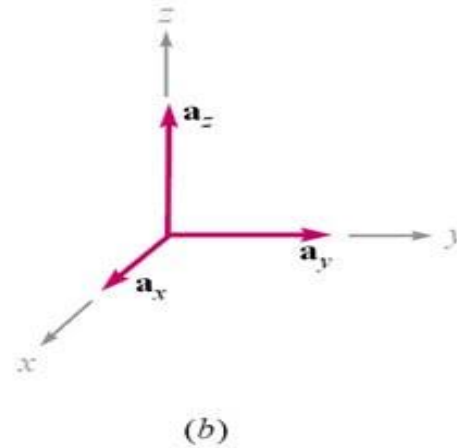
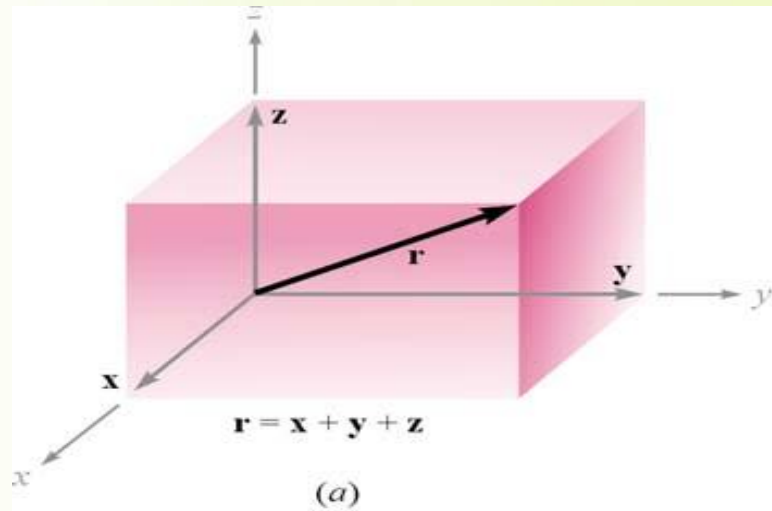
1. Cartesian or rectangular co-ordinate system
2. Cylindrical co-ordinate System
3. Spherical co-ordinate System



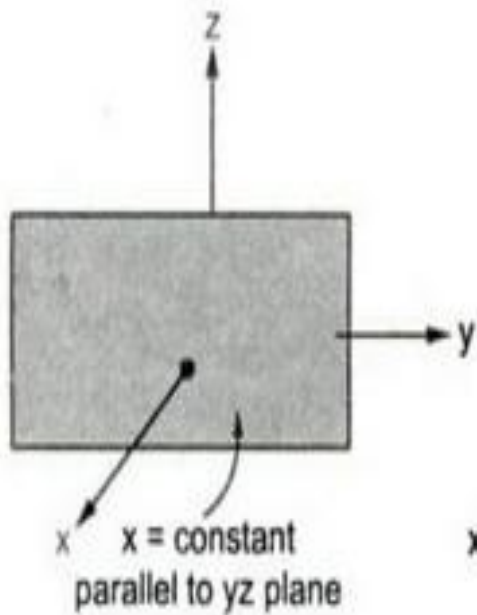
# The Cartesian Coordinate System : vertices are x,y,z



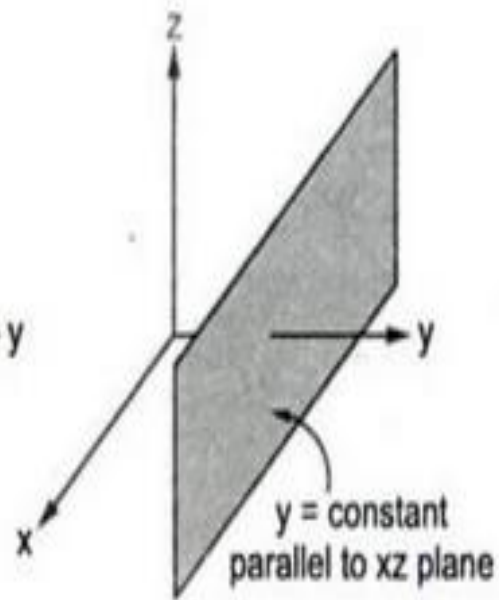
# Vector Components and Unit Vectors



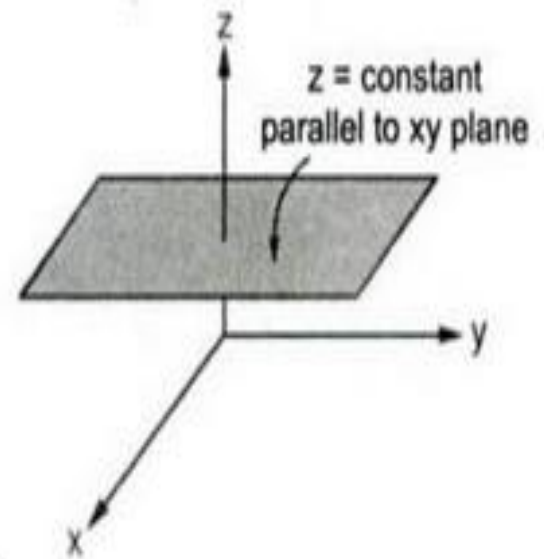
# Constant planes



(a)

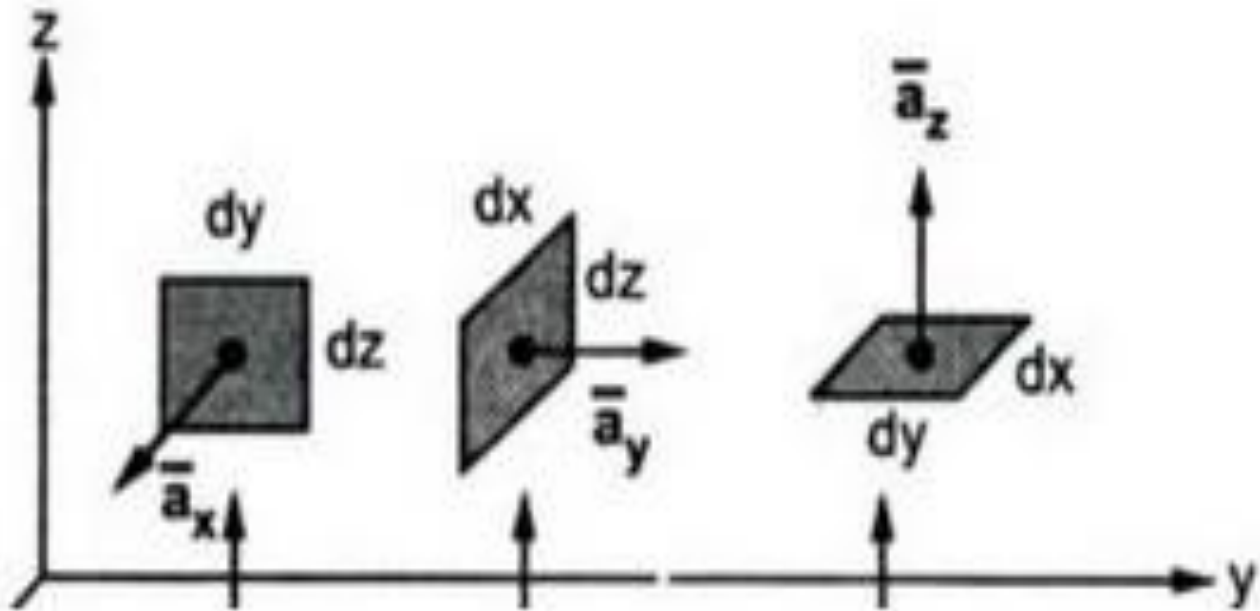


(b)



(c)

## Contd...

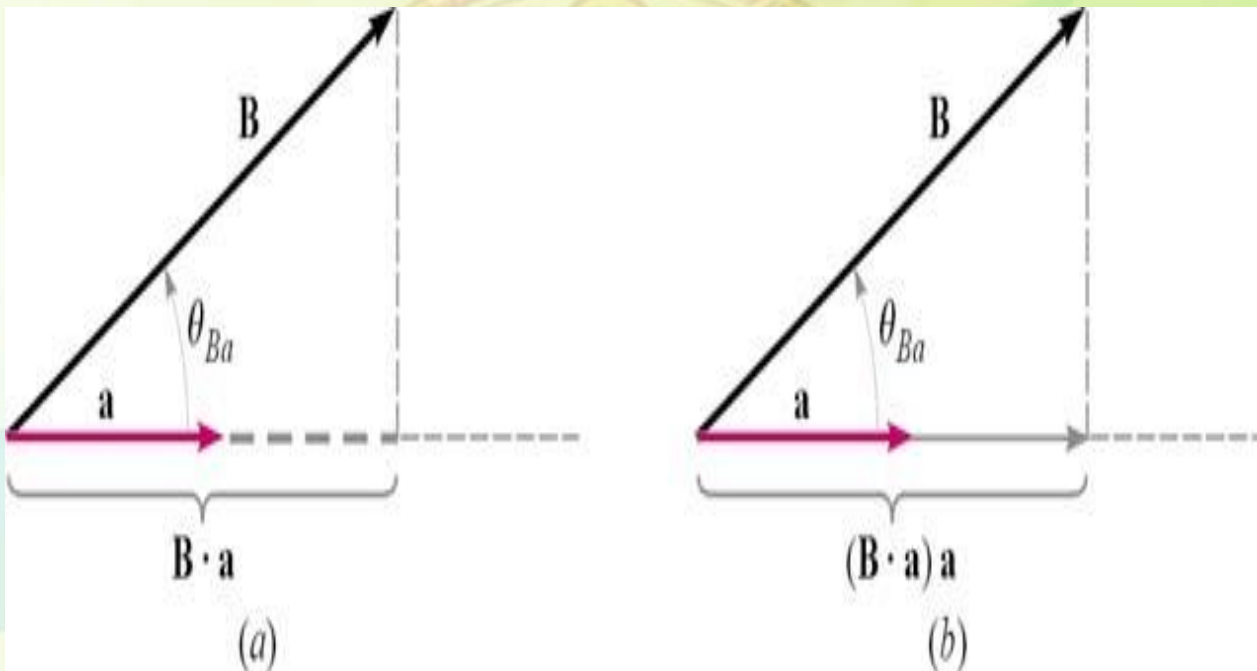


Distance vector -  $d\mathbf{l}$   
Differential surfaces -  
 $dv =$

# The Dot product

$$\mathbf{A} \cdot \mathbf{B} = |\mathbf{A}| |\mathbf{B}| \cos \theta_{AB}$$

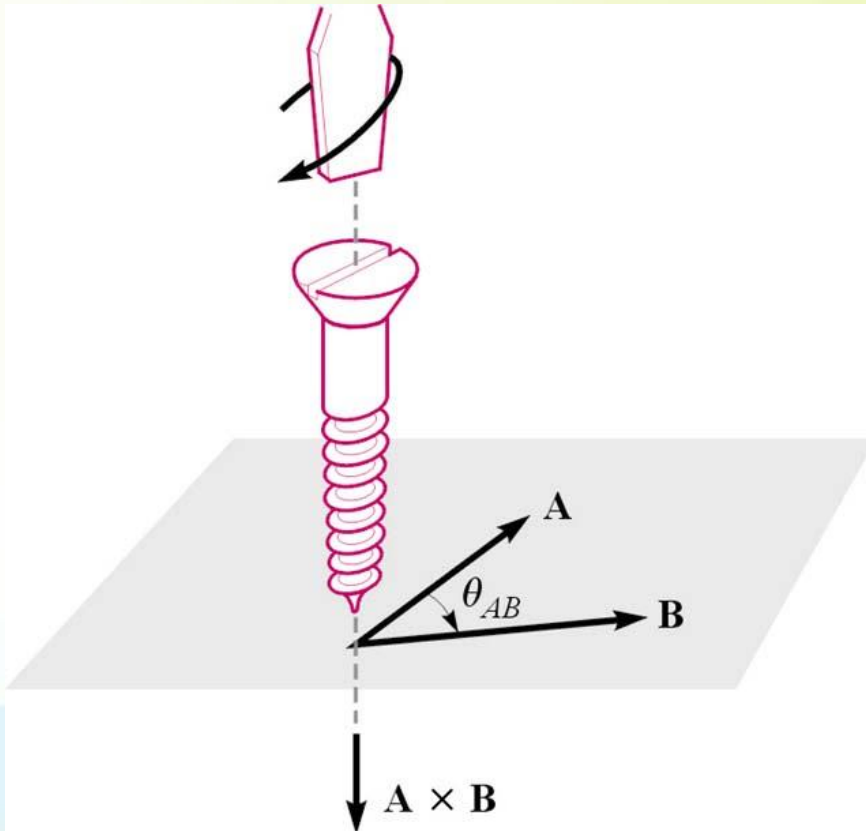
B in the direction of A  
You need to normalize a  
before the dot product.





# The Cross Product

$$\mathbf{A} \times \mathbf{B} = a_N |\mathbf{A}| |\mathbf{B}| \sin \theta_{AB}$$



$$\mathbf{A} \times \mathbf{B} = \begin{pmatrix} a_x & a_y & a_z \\ A_x & A_y & A_z \\ B_x & B_y & B_z \end{pmatrix}$$

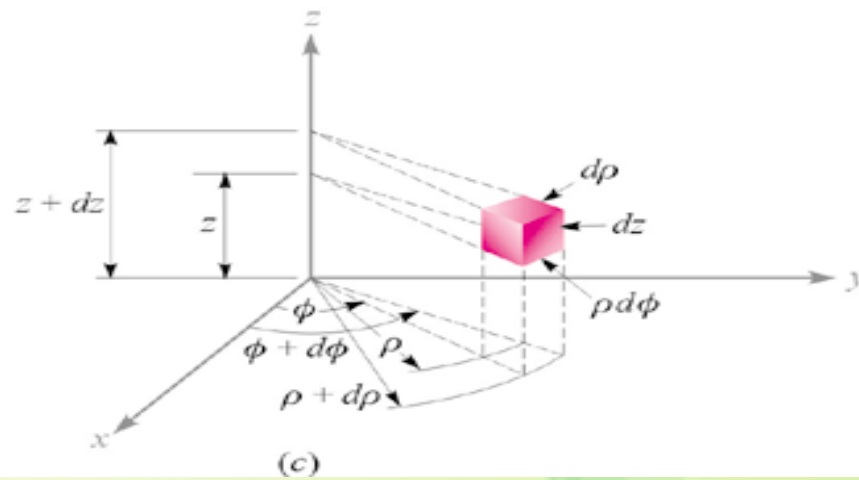
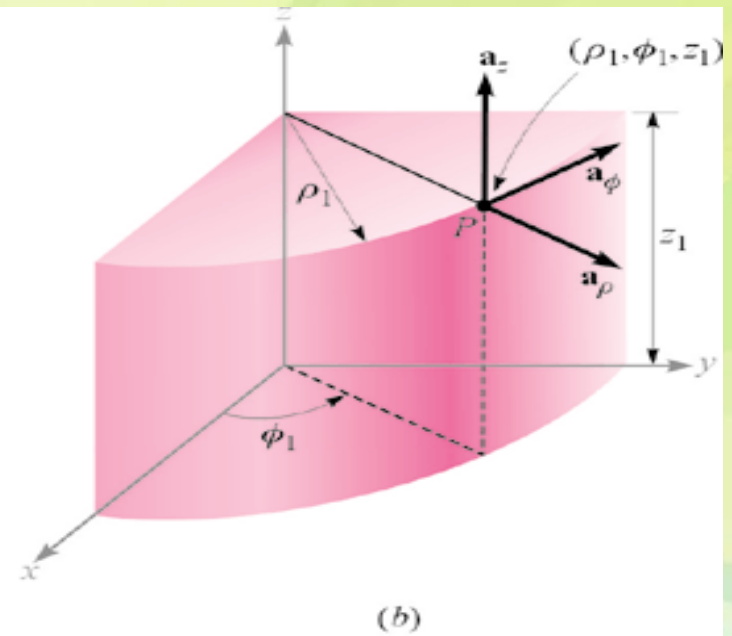
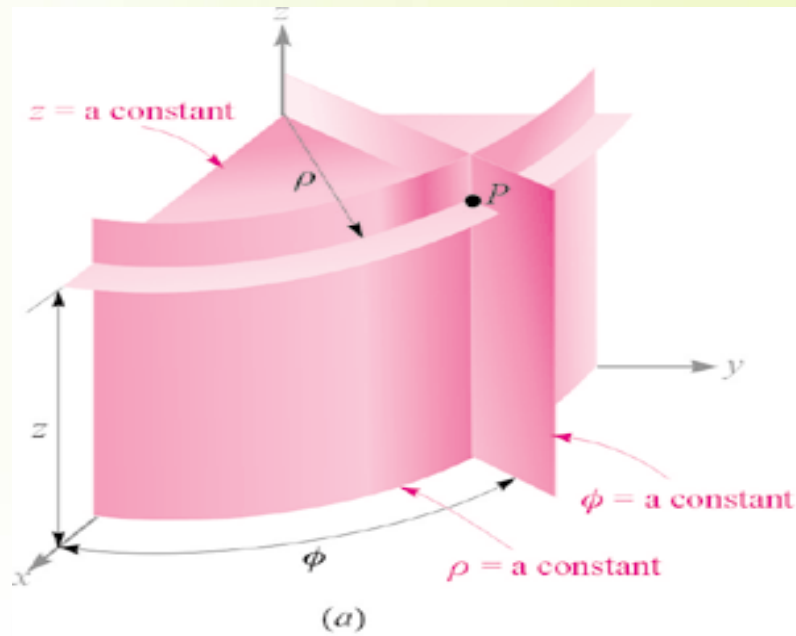
Example

$$\mathbf{A} := \begin{pmatrix} 2 \\ -3 \\ 1 \end{pmatrix}$$

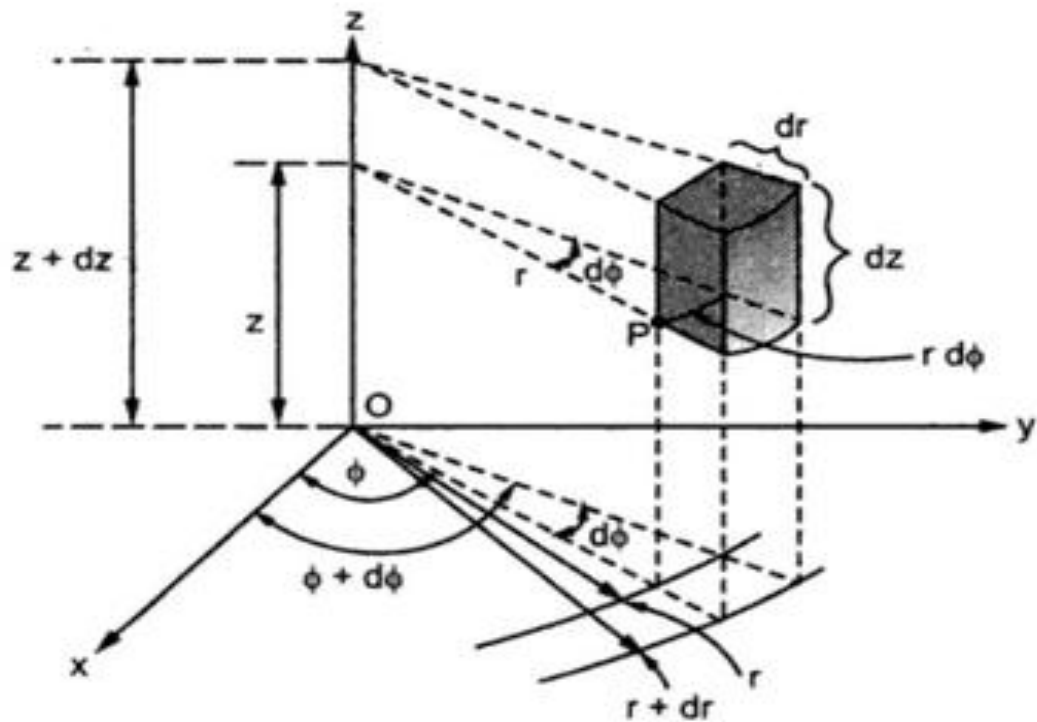
$$\mathbf{B} := \begin{pmatrix} -4 \\ -2 \\ 5 \end{pmatrix}$$

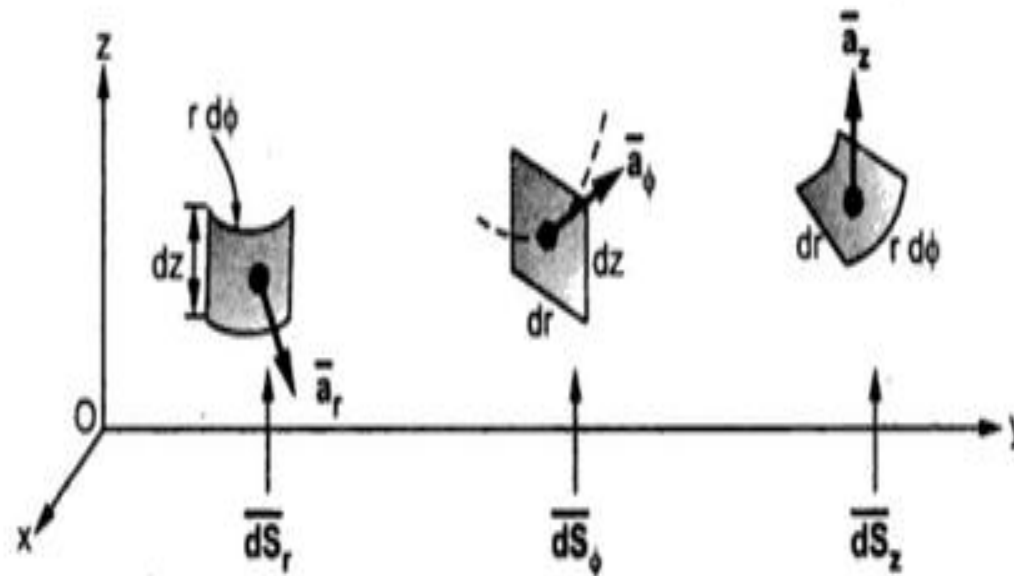
$$\mathbf{A} \times \mathbf{B} = \begin{pmatrix} -13 \\ -14 \\ -16 \end{pmatrix}$$

# Circular Cylindrical Coordinate System



# Differential volume

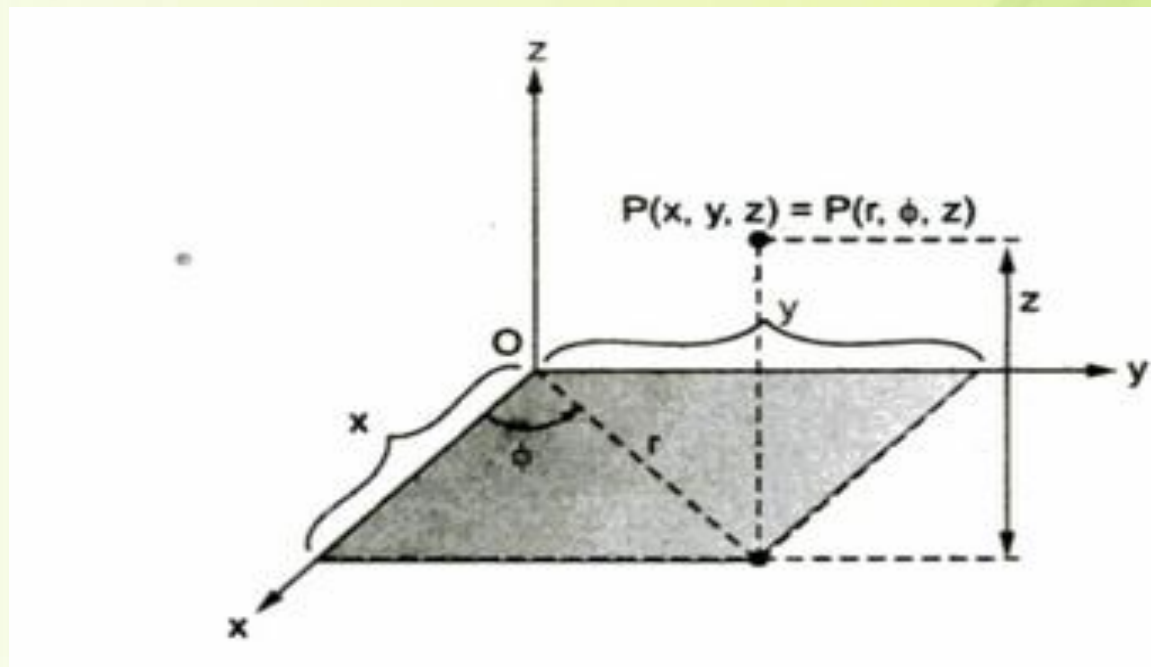




$\bar{dS}_r$  = Differential vector surface area normal to r direction  
 =  $r d\phi dz \bar{a}_r$

$\bar{dS}_\phi$  = Differential vector surface area normal to  $\phi$  direction  
 =  $dr dz \bar{a}_\phi$

$\bar{dS}_z$  = Differential vector surface area normal to z direction  
 =  $r dr d\phi \bar{a}_z$

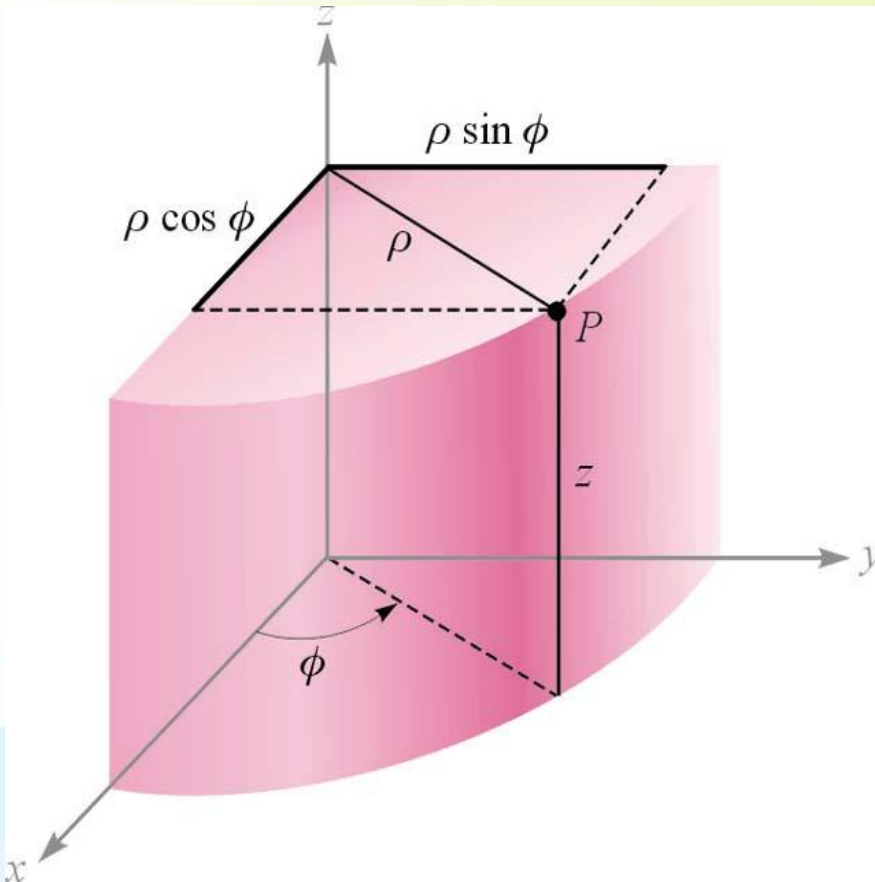


$$x = r \cos \phi \quad y = r \sin \phi \quad \text{and} \quad z = z$$

It can be seen that,  $r$  can be expressed in terms of  $x$  and  $y$  as,

$$r = \sqrt{x^2 + y^2}$$

# Circular Cylindrical Coordinate System



$$x = \rho \cdot \cos(\phi)$$

$$y = \rho \cdot \sin(\phi)$$

$$z = z$$

$$\rho = \sqrt{x^2 + y^2} \quad \rho \geq 0$$

$$\phi = \text{atan}\left(\frac{y}{x}\right)$$

$$z = z$$

## Dot Product

$$A = A_x \cdot a_x + A_y \cdot a_y + A_z \cdot a_z$$

$$A = A_\rho \cdot a_\rho + A_\phi \cdot a_\phi + A_z \cdot a_z$$

$$A_\rho = A \cdot a_\rho$$

$$A_\phi = A \cdot a_\phi$$

$$A_z = A_z$$

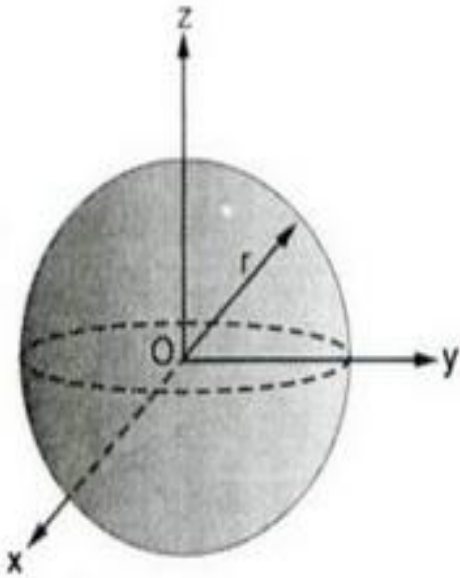
$$A_\rho = (A_x \cdot a_x + A_y \cdot a_y + A_z \cdot a_z) \cdot a_\rho = A_x \cdot a_x \cdot a_\rho + A_y \cdot a_y \cdot a_\rho$$

$$A_\phi = (A_x \cdot a_x + A_y \cdot a_y + A_z \cdot a_z) \cdot a_\phi = A_x \cdot a_x \cdot a_\phi + A_y \cdot a_y \cdot a_\phi$$

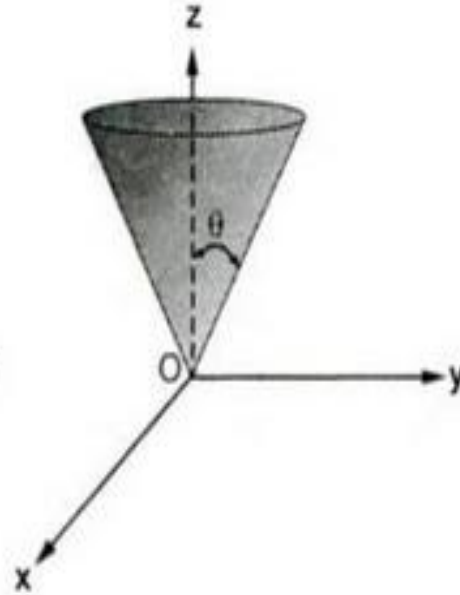
$$A_z = (A_x \cdot a_x + A_y \cdot a_y + A_z \cdot a_z) \cdot a_z = A_z \cdot a_z \cdot a_z = A_z$$

$$a_z \cdot a_\rho = a_z \cdot a_\phi = 0$$

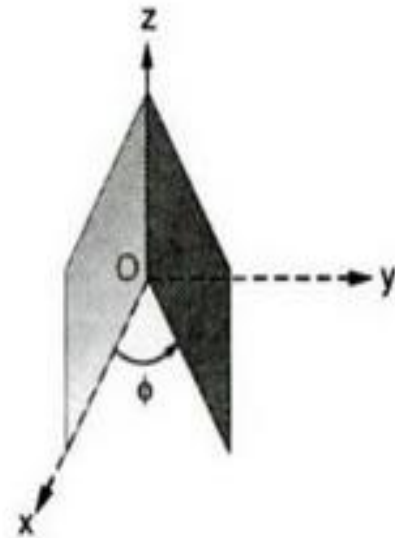
# Spherical co-ordinate system



(a) Sphere of radius  $r$  with centre as origin

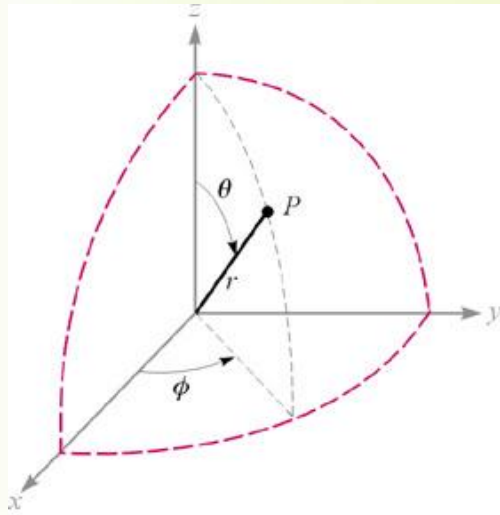


(b) Right circular cone with apex at origin

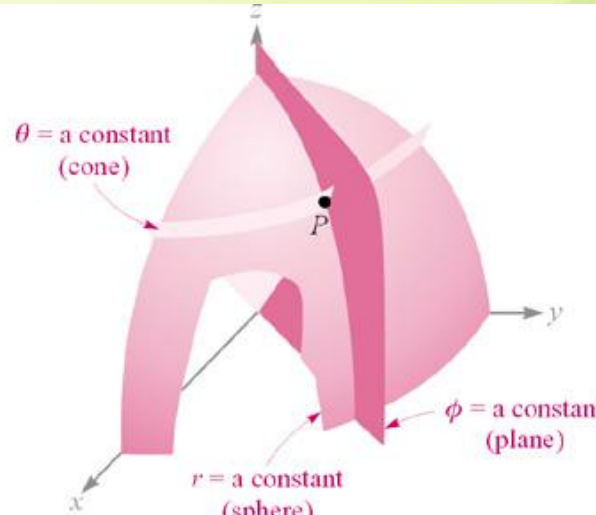


(c) Half plane perpendicular to  $xy$  plane

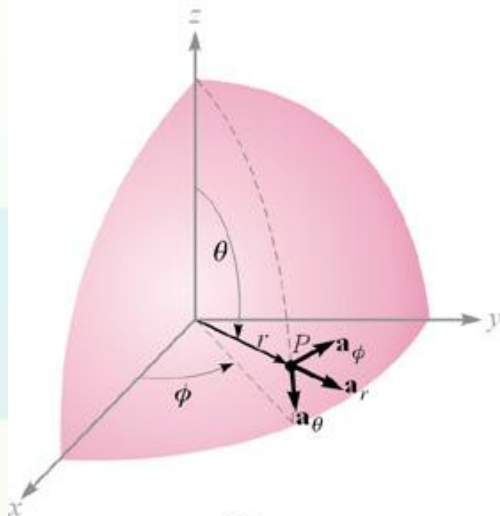
# The Spherical Coordinate System



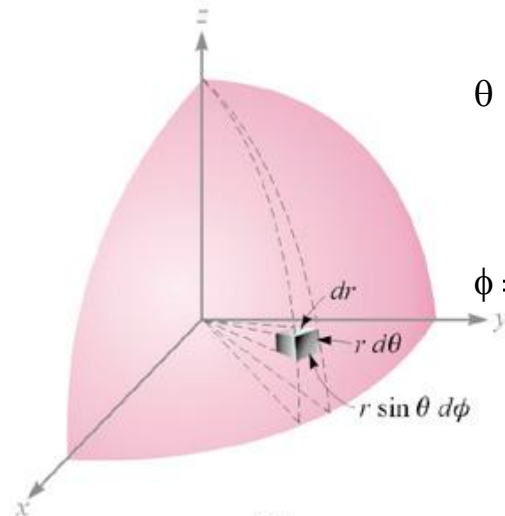
(a)



(b)



(c)



(d)

$$x = r \cdot \sin(\theta) \cdot \cos(\phi)$$

$$y = r \cdot \sin(\theta) \cdot \sin(\phi)$$

$$z = r \cdot \cos(\theta)$$

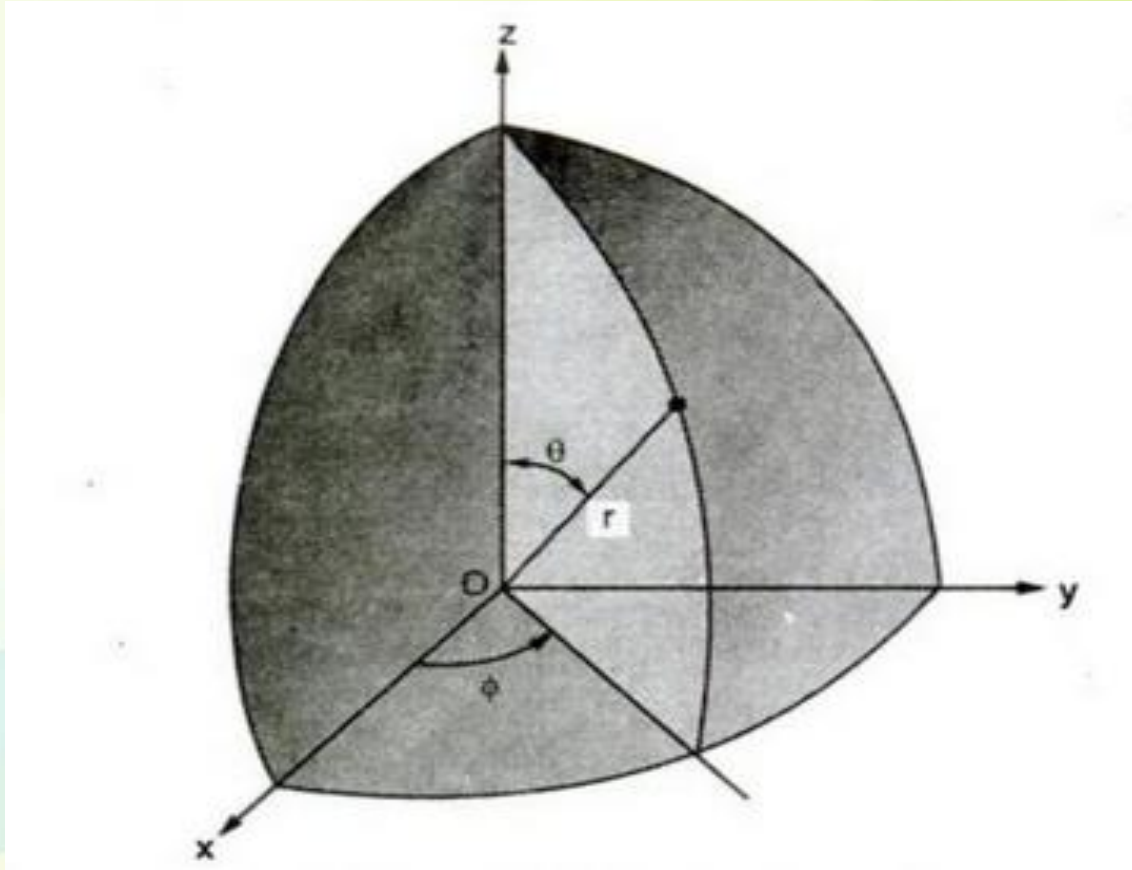
$$r = \sqrt{x^2 + y^2 + z^2} \quad r \geq 0$$

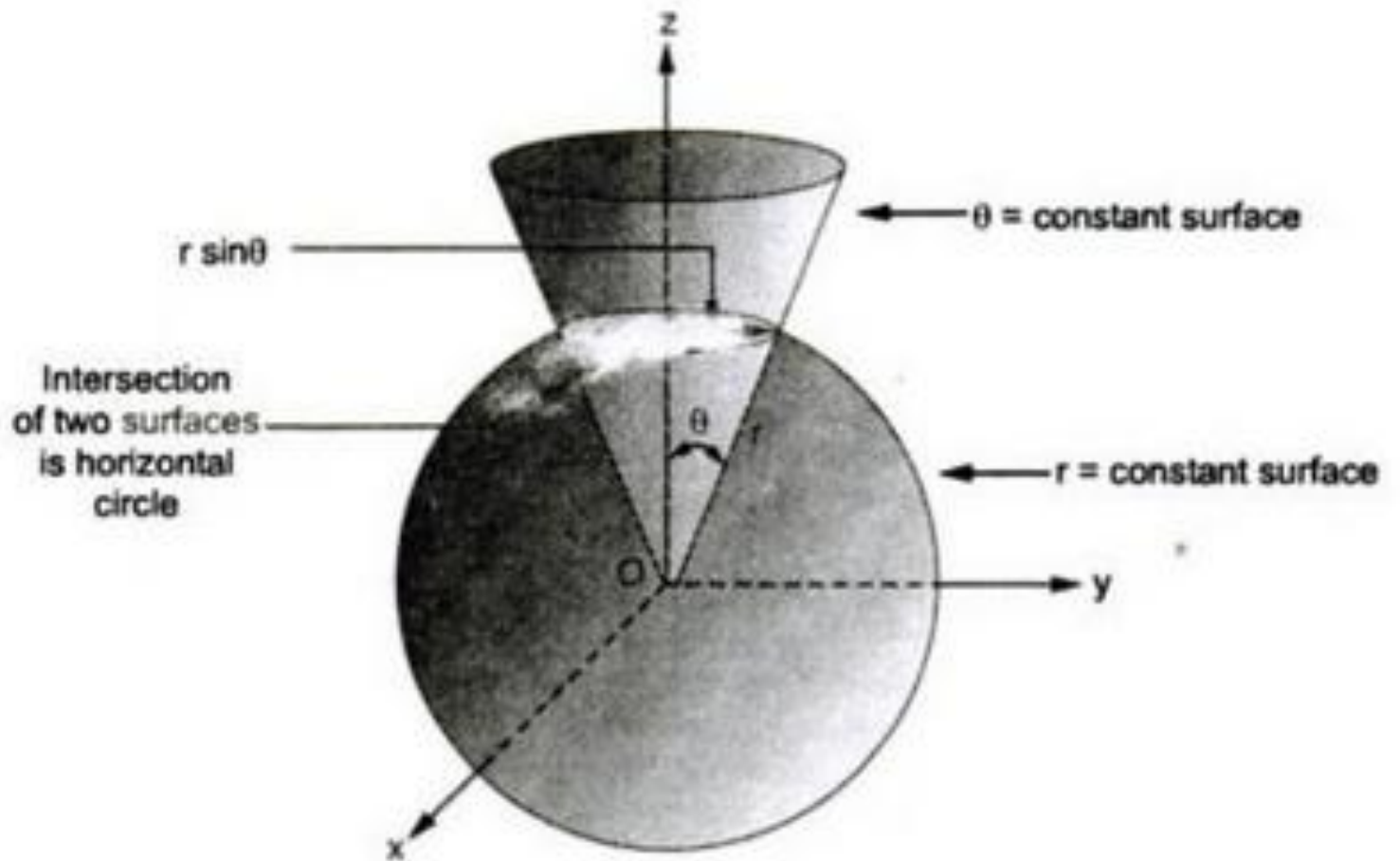
$$\theta = \arccos\left(\frac{z}{\sqrt{x^2 + y^2 + z^2}}\right) \quad 0 \leq \theta \leq 180^\circ$$

$$\phi = \operatorname{atan}\left(\frac{y}{x}\right)$$

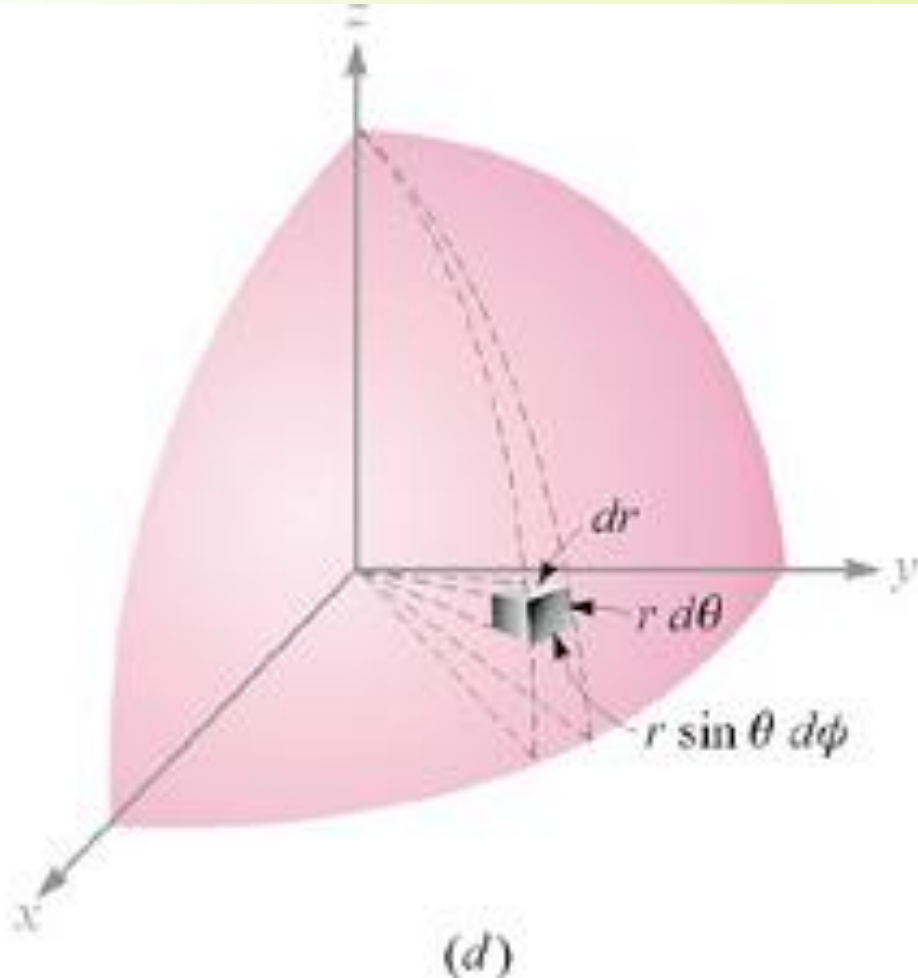


# Spherical Co-ordinate System





# The Spherical Coordinate System



$$x = r \cdot \sin(\theta) \cdot \cos(\phi)$$

$$y = r \cdot \sin(\theta) \cdot \sin(\phi)$$

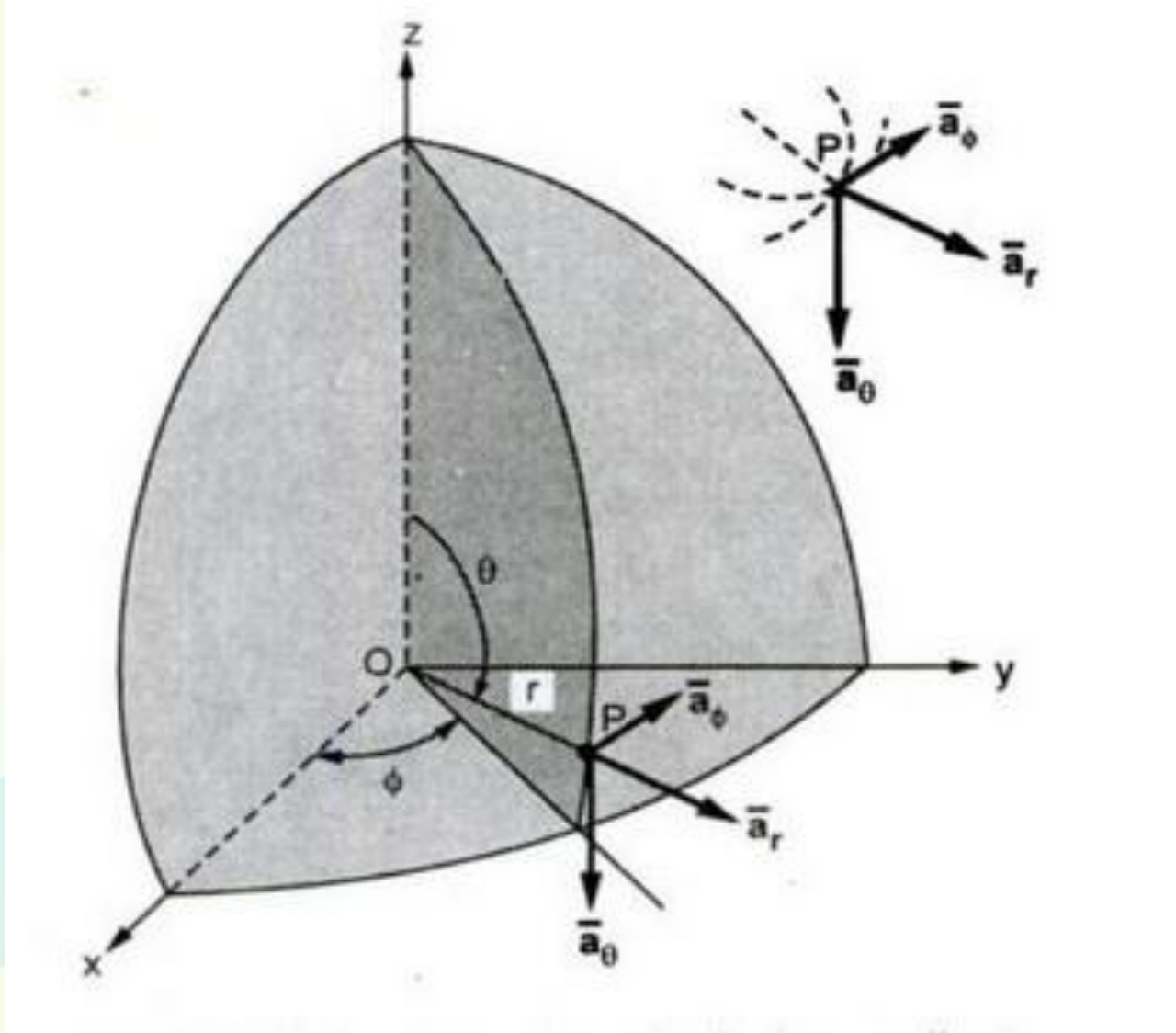
$$z = r \cdot \cos(\theta)$$

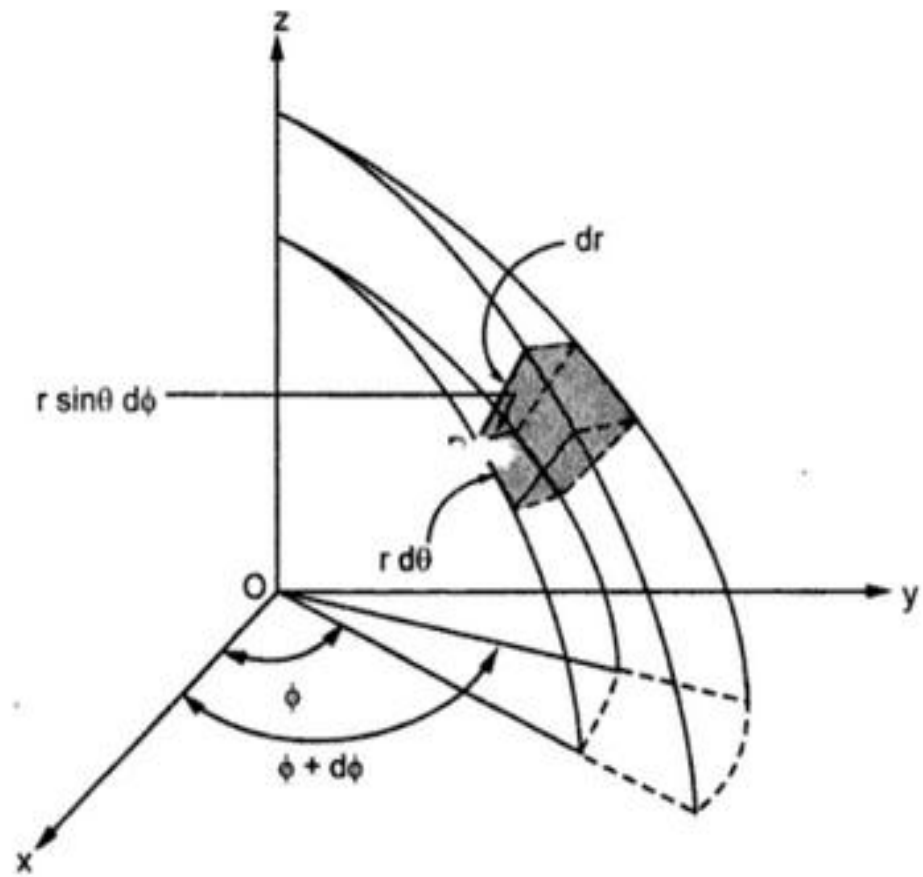
$$r \cdot dr \cdot d\theta$$

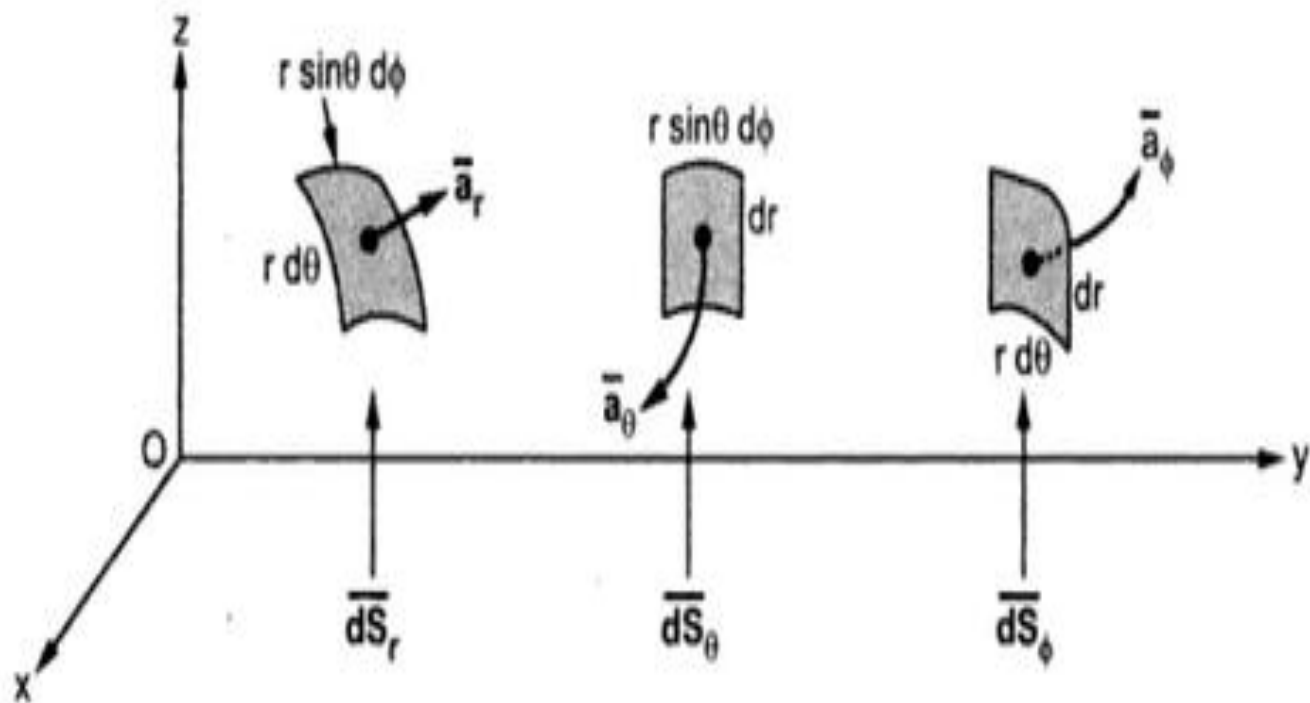
$$r \cdot \sin(\theta) \cdot dr \cdot d\phi$$

$$r^2 \cdot \sin(\theta) \cdot d\theta \cdot d\phi$$

$$r^2 \cdot \sin(\theta) \cdot dr \cdot d\theta \cdot d\phi$$



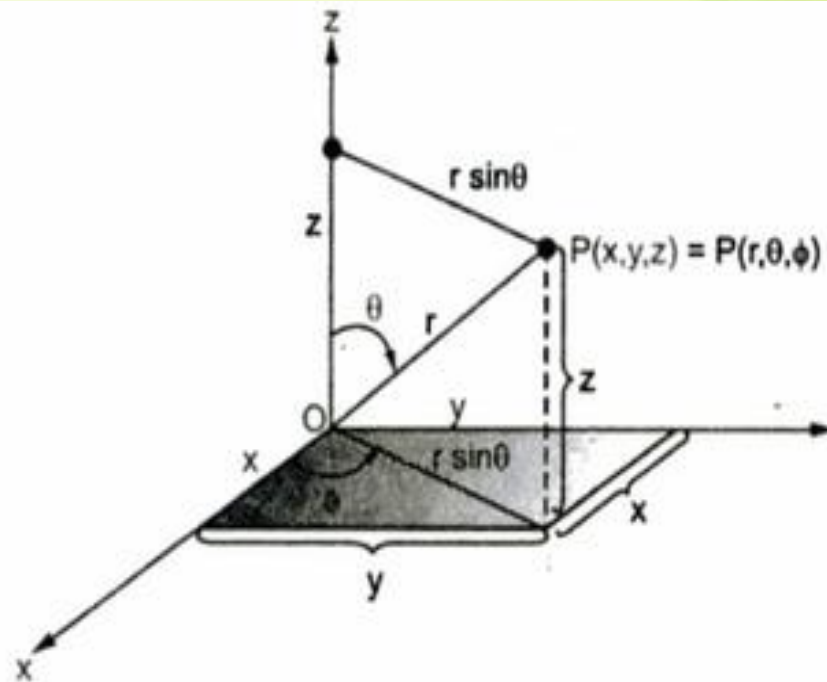




$\bar{dS}_r$  = Differential vector surface area normal to  $r$  direction  
 =  $r^2 \sin \theta d\theta d\phi$

$\bar{dS}_\theta$  = Differential vector surface area normal to  $\theta$  direction  
 =  $r \sin \theta dr d\phi$

$\bar{dS}_\phi$  = Differential vector surface area normal to  $\phi$  direction  
 =  $r dr d\theta$



$$x = r \sin \theta \cos \phi \quad \text{and} \quad y = r \sin \theta \sin \phi$$

$$z = r \cos \theta$$

$$x = r \sin \theta \cos \phi, \quad y = r \sin \theta \sin \phi \quad \text{and} \quad z = r \cos \theta$$

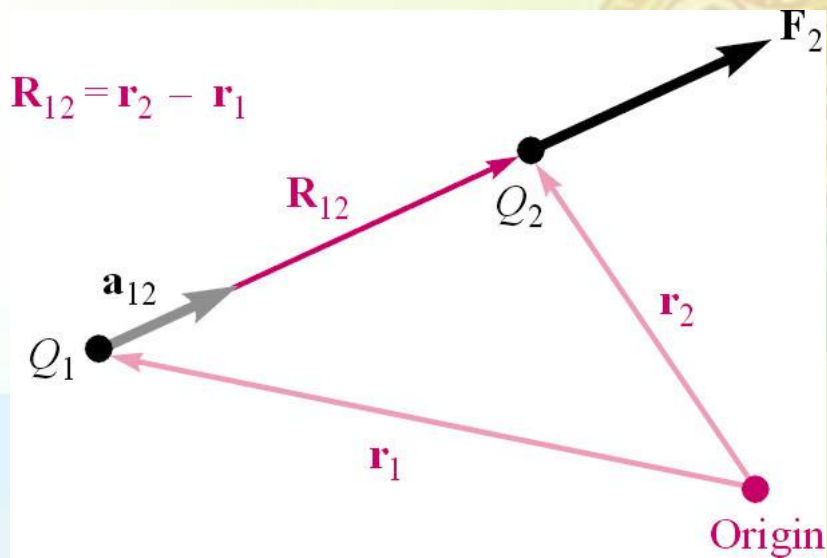
# The Experimental Law of Coulomb

$$F = k \cdot \frac{Q_1 \cdot Q_2}{R^2}$$

$$k = \frac{1}{4 \cdot \pi \cdot \epsilon_0}$$

$$\epsilon_0 = 8.85410^{-12} = \frac{1}{36\pi} \cdot 10^{-9} \quad \frac{\text{F}}{\text{m}}$$

$$\longrightarrow F = \frac{Q_1 \cdot Q_2}{4 \cdot \pi \cdot \epsilon_0 \cdot R^2}$$



$$F = \frac{Q_1 \cdot Q_2}{4 \cdot \pi \cdot \epsilon_0 \cdot R^2} \cdot a_{12}$$

$$a_{12} = \frac{\mathbf{r}_2 - \mathbf{r}_1}{|\mathbf{r}_2 - \mathbf{r}_1|}$$

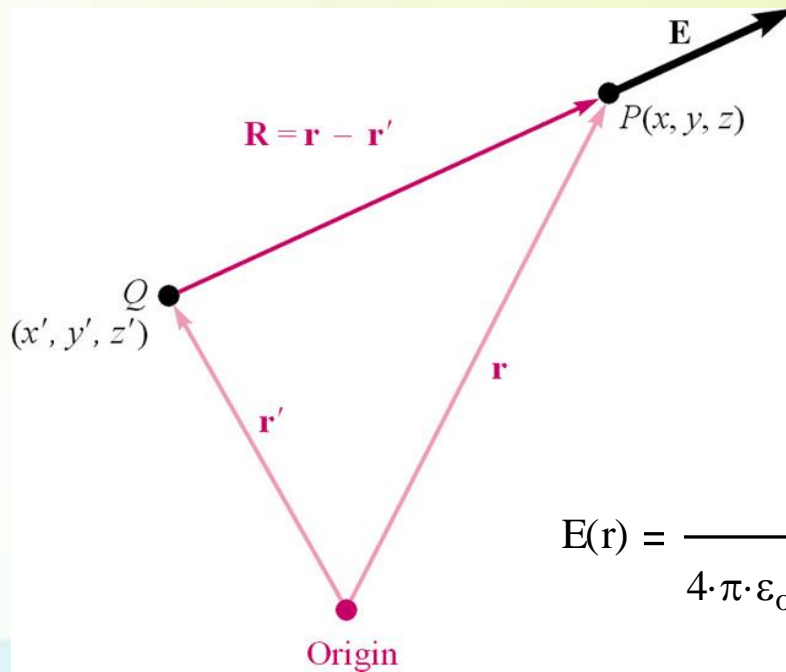


# Electric Field Intensity

$$F_t = \frac{Q_1 \cdot Q_t}{4 \cdot \pi \cdot \epsilon_0 \cdot (R_{1t})^2} \cdot a_{1t}$$

$$\frac{F_t}{Q_t} = \frac{Q_1}{4 \cdot \pi \cdot \epsilon_0 \cdot (R_{1t})^2} \cdot a_{1t}$$

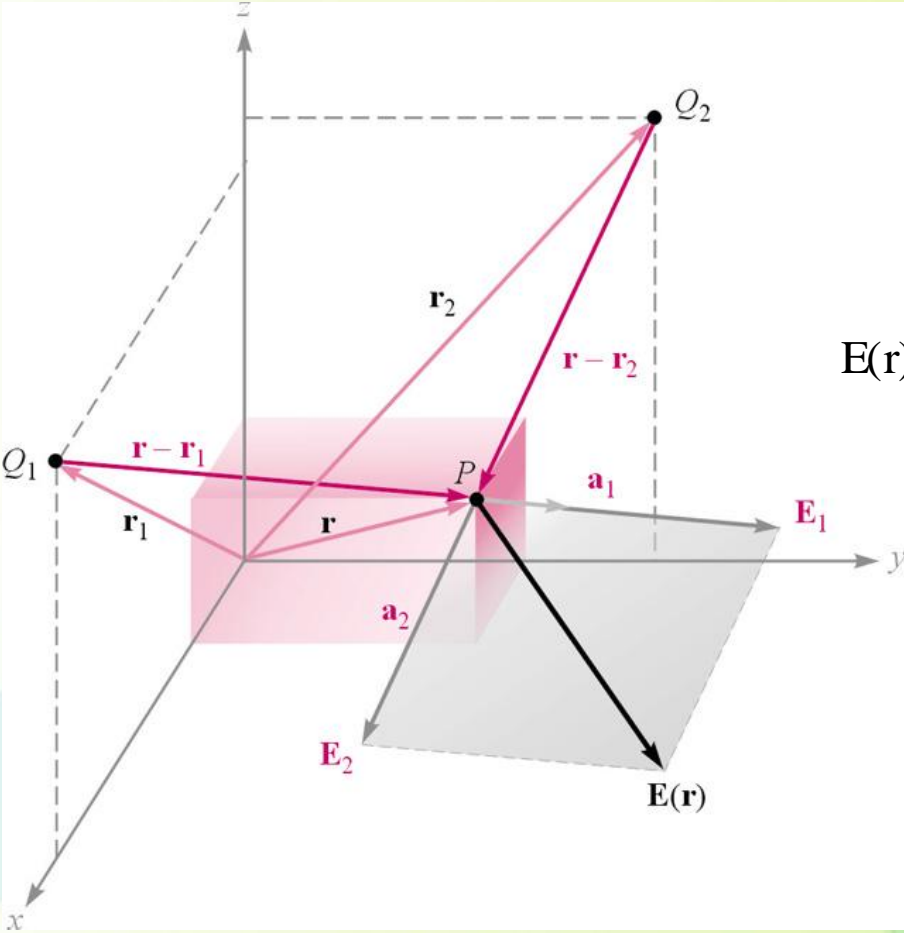
$$E = \frac{F_t}{Q_t}$$



$$E(\mathbf{r}) = \frac{Q}{4 \cdot \pi \cdot \epsilon_0 \cdot (|\mathbf{r} - \mathbf{r}_1|)^2} \cdot \frac{\mathbf{r} - \mathbf{r}_1}{|\mathbf{r} - \mathbf{r}_1|}$$

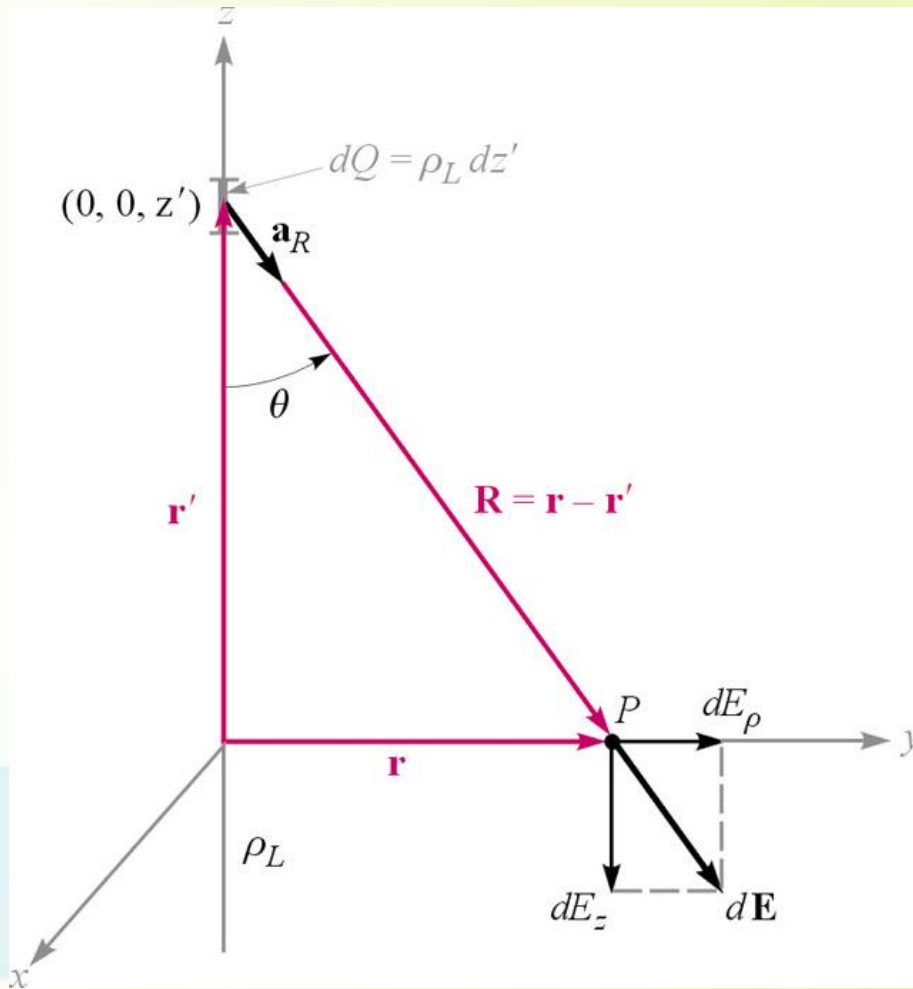
$$E(\mathbf{r}) = \frac{Q \cdot [(x - x_1) \cdot \mathbf{a}_x + (y - y_1) \cdot \mathbf{a}_y + (z - z_1) \cdot \mathbf{a}_z]}{4 \cdot \pi \cdot \epsilon_0 \cdot [(x - x_1)^2 + (y - y_1)^2 + (z - z_1)^2]^{\frac{3}{2}}}$$

# Electric Field Intensity



$$\mathbf{E}(\mathbf{r}) = \sum_{m=1}^n \frac{Q_m}{4 \cdot \pi \cdot \epsilon_0 \cdot (|\mathbf{r} - \mathbf{r}_m|)^2} \cdot \mathbf{a}_m$$

# Field of a Line Charge



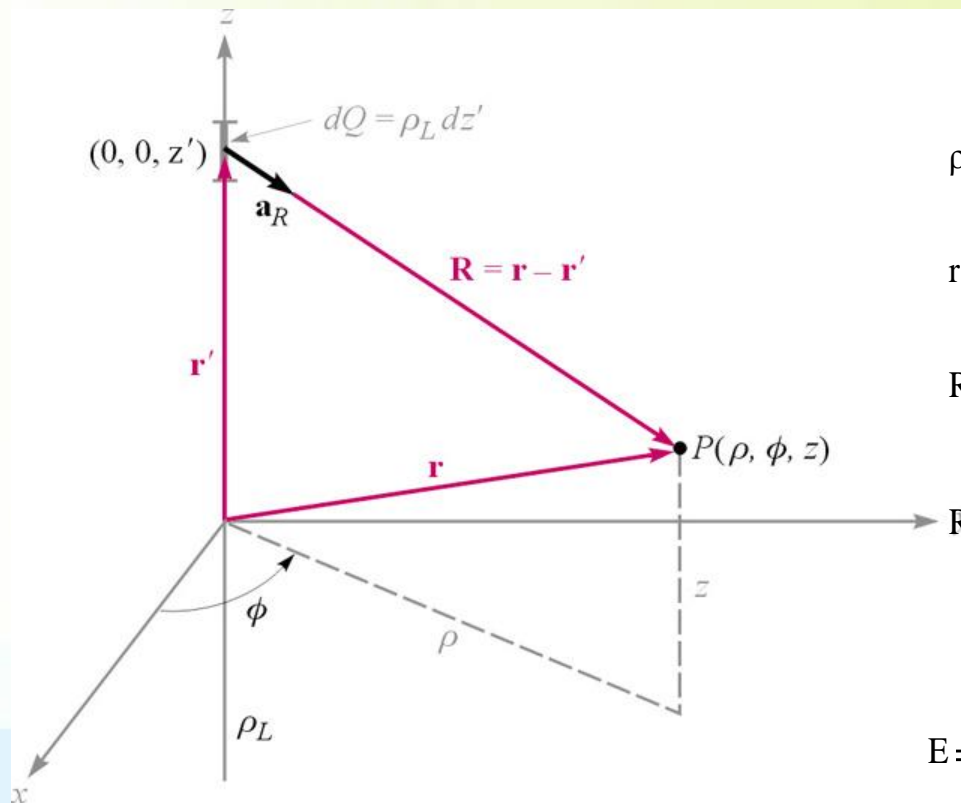
$$E_\rho = \int_{-\Omega}^{\Omega} \frac{\rho_L \cdot \rho}{4 \cdot \pi \cdot \epsilon_0 \cdot (\rho^2 + z^2)^{\frac{3}{2}}} dz$$

$$E_\rho = \frac{\rho_L}{2 \cdot \pi \cdot \epsilon_0 \cdot \rho}$$

$$\mathbf{E} = \frac{\rho_L}{2 \cdot \pi \cdot \epsilon_0 \cdot \rho} \cdot \mathbf{a}_\rho$$

# Field of a Line Charge (neglect symmetry)

$$E = \int \int \int \frac{\rho v q}{4 \cdot \pi \cdot \epsilon_0} \cdot \frac{(r - r_1)}{(|r - r_1|)^3} dx_1 dy_1 dz_1$$



$$\rho v_1 = \rho_L \cdot dz_1$$

$$r = \rho \cdot a_\rho + z \cdot a_z \quad r_1 = z_1 \cdot a_z$$

$$R = r - r_1 = \rho \cdot a_\rho + (z - z_1) \cdot a_z$$

$$R = \sqrt{\rho^2 + (z - z_1)^2}$$

$$a_R = \frac{\rho \cdot a_\rho + (z - z_1) \cdot a_z}{\sqrt{\rho^2 + (z - z_1)^2}}$$

$$E = \int_{-\Omega}^{\Omega} \frac{(\rho_L \cdot dz_1) \cdot [\rho \cdot a_\rho + (z - z_1) \cdot a_z]}{[4 \cdot \pi \cdot \epsilon_0 \cdot [\rho^2 + (z - z_1)^2]]^{\frac{3}{2}}} dz_1$$

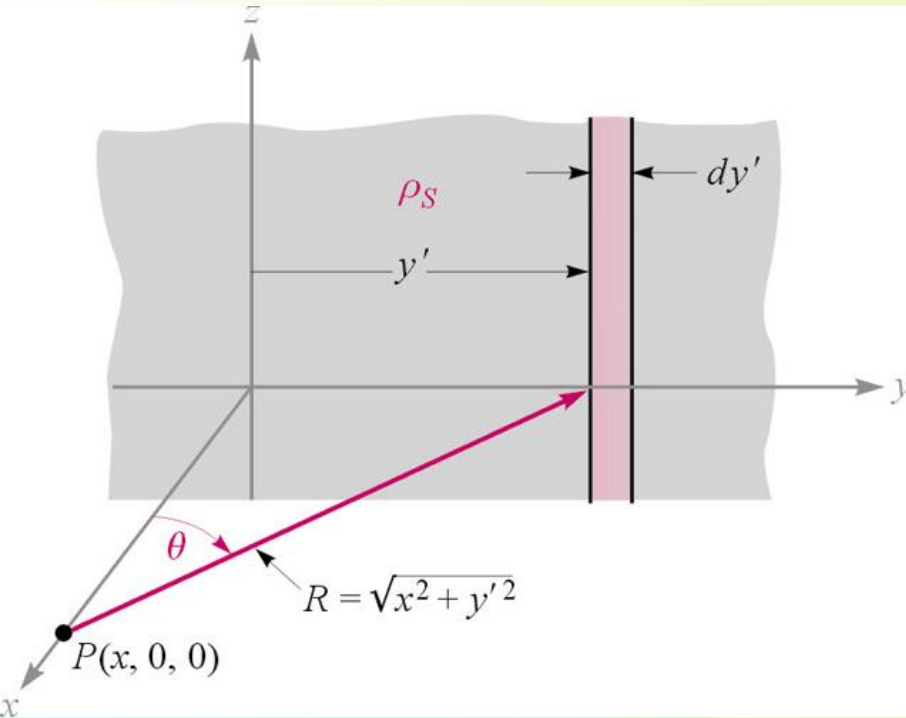
# Field of a Line Charge (neglect symmetry)

$$E = \frac{\rho L}{(4 \cdot \pi \cdot \epsilon_0)} \cdot \left[ a_\rho \cdot \int_{-\Omega}^{\Omega} \frac{(\rho \cdot dz_1)}{\left[ \left[ \rho^2 + (z - z_1)^2 \right] \right]^{\frac{3}{2}}} dz_1 + a_z \cdot \int_{-\Omega}^{\Omega} \frac{((z - z_1))}{\left[ \left[ \rho^2 + (z - z_1)^2 \right] \right]^{\frac{3}{2}}} dz_1 \right]$$

$$E = \frac{\rho L}{(4 \cdot \pi \cdot \epsilon_0)} \cdot \left[ a_\rho \cdot \rho \cdot \frac{1}{\rho^2} \cdot \frac{-\overset{-\Omega \text{ to } \Omega}{(z - z_1)}}{\sqrt{\rho^2 + (z - z_1)^2}} + a_z \cdot \frac{1}{\sqrt{\rho^2 + (z - z_1)^2}} \right]$$

$$E = \frac{\rho L}{(4 \cdot \pi \cdot \epsilon_0)} \cdot \left( a_\rho \cdot \frac{2}{\rho} + a_z \cdot 0 \right) = \frac{\rho L}{(2 \cdot \pi \cdot \epsilon_0) \cdot \rho} \cdot a_\rho$$

# Field of a Sheet of Charge



$$R = \sqrt{x^2 + y'^2}$$

$$dE_x = \rho_s \cdot \frac{dy'}{2 \cdot \pi \cdot \epsilon_0 \cdot \sqrt{x^2 + y'^2}} \cdot \cos(\theta) = \frac{\rho_s}{(2 \cdot \pi \cdot \epsilon_0)} \cdot \frac{x \cdot dy'}{x^2 + y'^2}$$

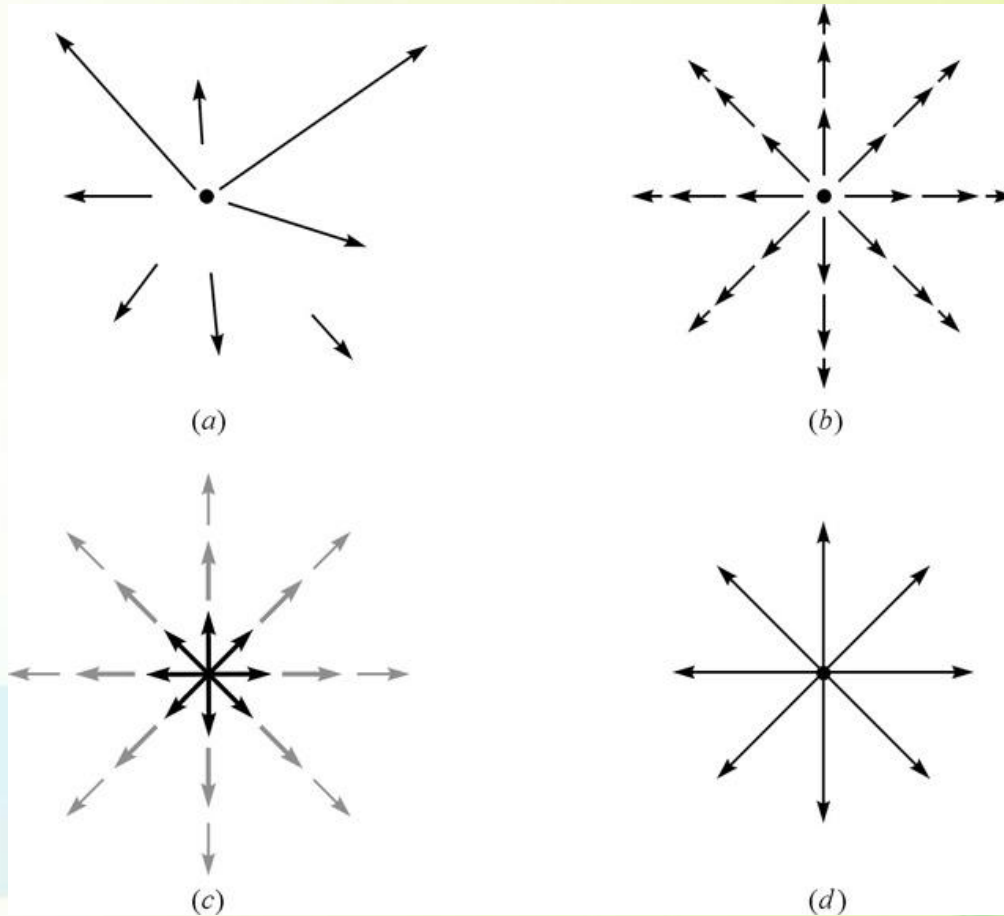
$$E_x = \frac{\rho_s}{(2 \cdot \pi \cdot \epsilon_0)} \cdot \int_{-\Omega}^{\Omega} \frac{x}{x^2 + y'^2} dy' = \frac{\rho_s}{(2 \cdot \pi \cdot \epsilon_0)} \cdot \text{atan}\left(\frac{y'}{x}\right) \Big|_{-\Omega}^{\Omega}$$

$$E_x = \frac{\rho_s}{2 \cdot \epsilon_0}$$

$$E = \frac{\rho_s}{2 \cdot \epsilon_0} \cdot a_n$$

This is a very interesting result. The field is constant in magnitude and direction. It is as strong a million miles away from the sheet as it is right of the surface.

# Streamlines and Sketches of Fields

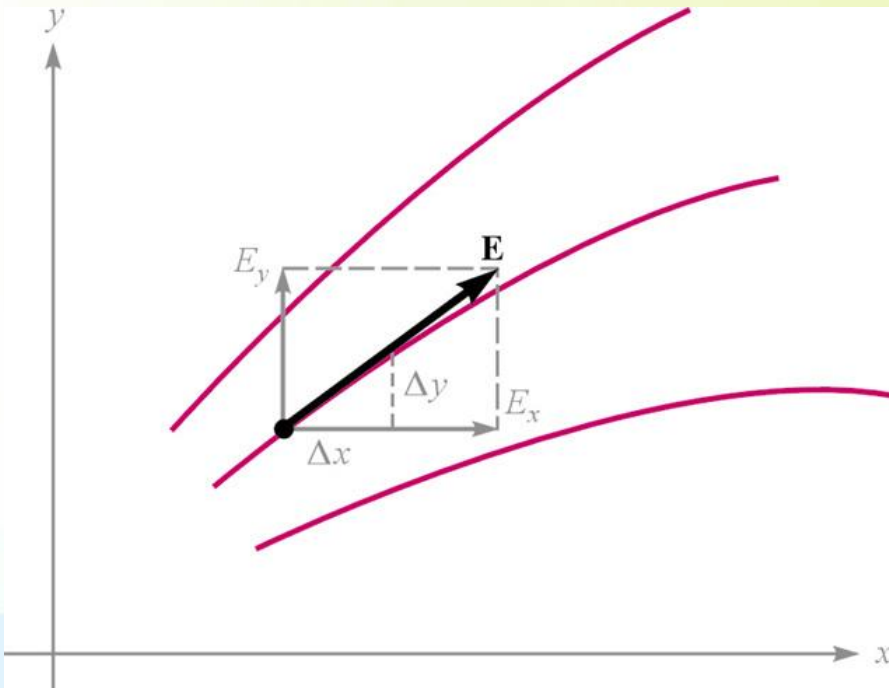


Cross-sectional view of the line charge.

Lengths proportional to the magnitudes of  $E$  and pointing in the direction of  $E$

# Streamlines and Sketches of Fields

$$\frac{E_y}{E_x} = \frac{dy}{dx}$$



$$\mathbf{E} = \frac{1}{\rho} \cdot \mathbf{a}_\rho$$

$$\mathbf{E} = \frac{x}{x^2 + y^2} \cdot \mathbf{a}_x + \frac{y}{x^2 + y^2} \cdot \mathbf{a}_y$$

$$\frac{dy}{dx} = \frac{E_y}{E_x} = \frac{y}{x}$$

$$\frac{dy}{y} = \frac{dx}{x}$$

$$\ln(y) = \ln(x) + C_1$$

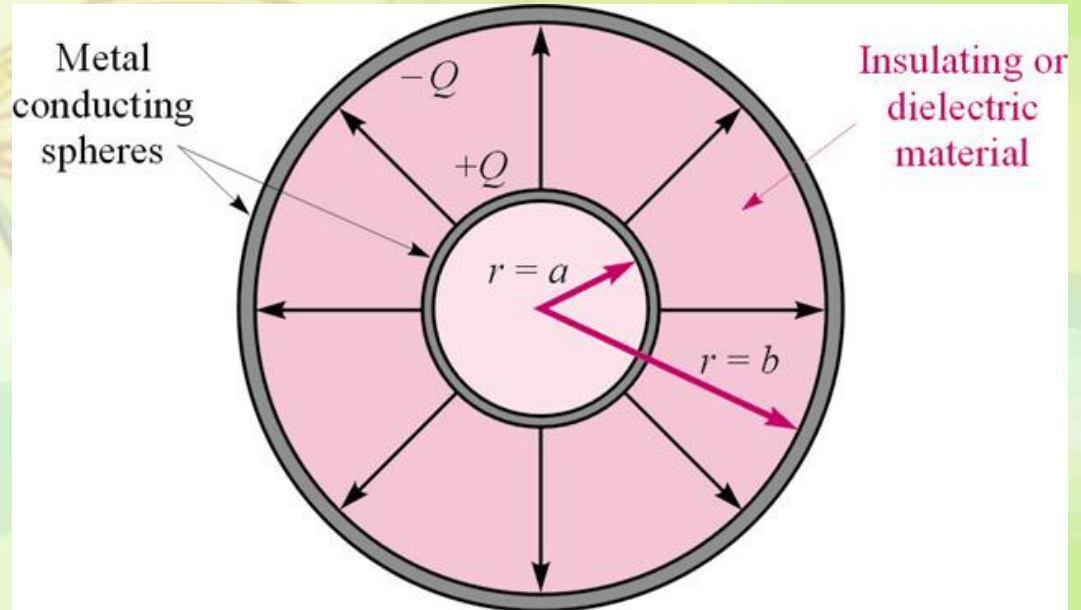
$$\ln(y) = \ln(x) + \ln(c)$$

$$y = C \cdot x$$



# 3.1 Electric Flux Density

- Faraday's Experiment



Flux =  $\Psi$ , same units as  $Q$

$\Psi$  is responsible for creating  $-Q$  on outer sphere

# Electric Flux Density, $\mathbf{D}$

- Units:  $\text{C/m}^2$
- Magnitude: Number of flux lines (coulombs) crossing a surface normal to the lines divided by the surface area.
- Direction: Direction of flux lines (same direction as  $\mathbf{E}$ ).
- For a point charge:  $\mathbf{D} = \frac{Q}{4\pi r^2} \mathbf{a}_r$
- For a general charge distribution,

$$\mathbf{D} = \epsilon_0 \mathbf{E} = \int_{\text{vol}} \frac{\rho_v d\mathbf{v}}{4\pi R^2} \mathbf{a}_r$$

### D3.1

Given a 60- $\mu\text{C}$  point charge located at the origin, find the total electric flux passing through:

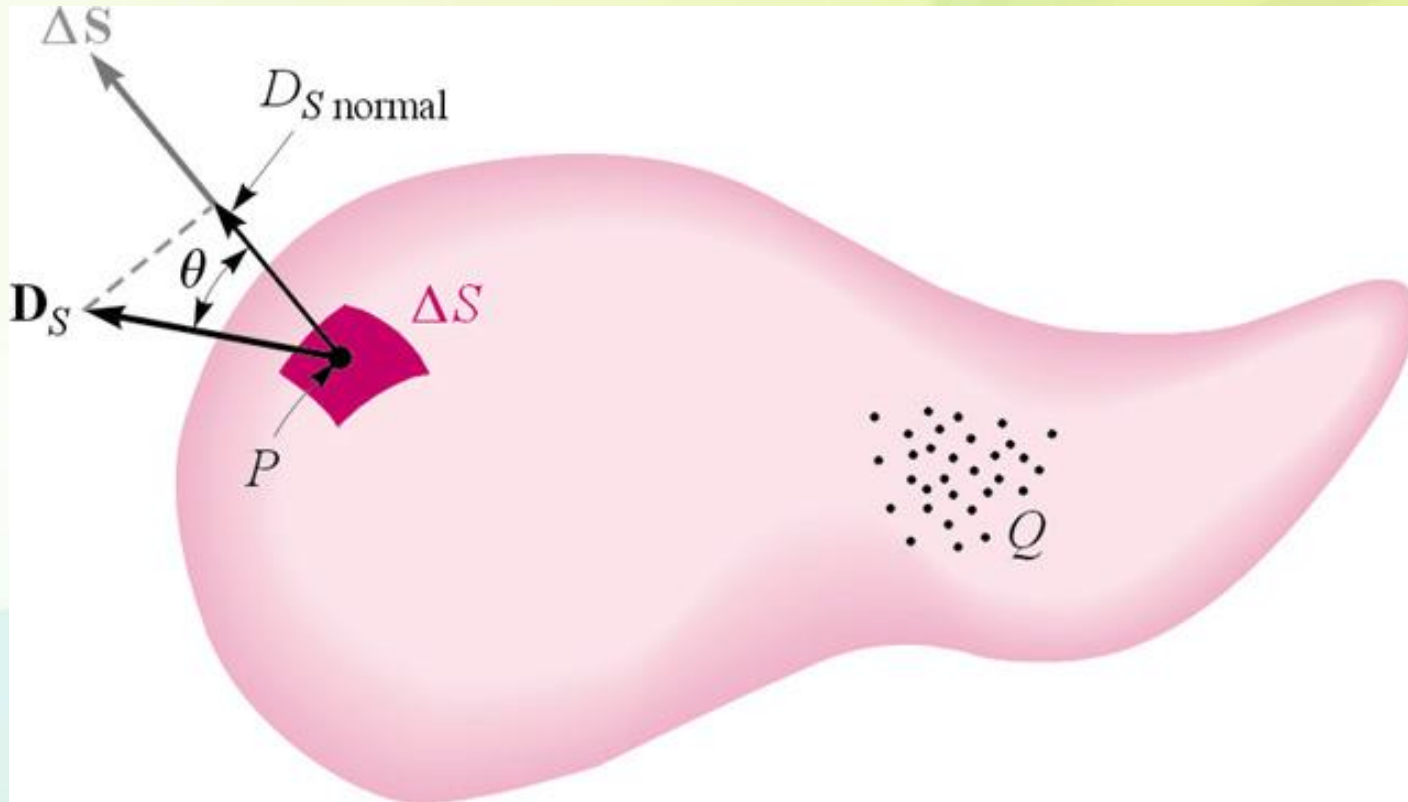
- (a) That portion of the sphere  $r = 26$  cm bounded by  $0 < \theta < \pi/2$  and  $0 < \phi < \pi/2$

# Gauss's Law

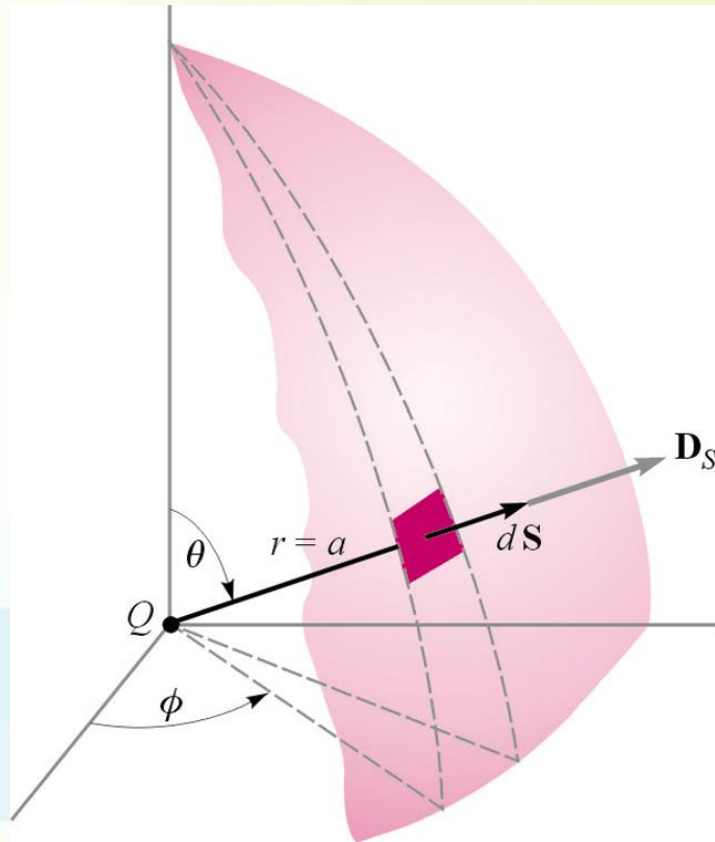
- “The electric flux passing through any closed surface is equal to the total charge enclosed by that surface.”

$$\Psi = \oint_S \mathbf{D}_S \cdot d\mathbf{S} = \text{charge enclosed} = Q$$

- The integration is performed over a *closed* surface, i.e. *gaussian surface*.



- We can check Gauss's law with a point charge example.



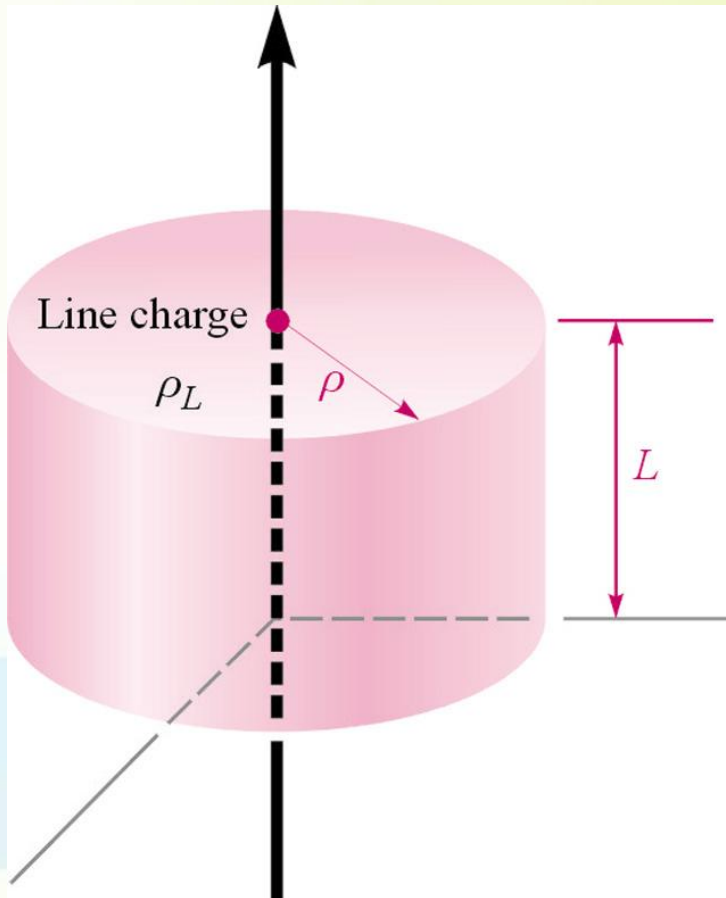
$$\int_0^{2\pi} \int_0^{\pi} \frac{q}{4\pi a^2} a^2 \sin[\theta] d\theta d\phi$$

q

# Symmetrical Charge Distributions

- Gauss's law is useful under two conditions.
  1.  $\mathbf{D}_S$  is everywhere either normal or tangential to the closed surface, so that  $\mathbf{D}_S \cdot d\mathbf{S}$  becomes either  $D_S dS$  or zero, respectively.
  1. On that portion of the closed surface for which  $\mathbf{D}_S \cdot d\mathbf{S}$  is not zero,  $D_S = \text{constant}$ .

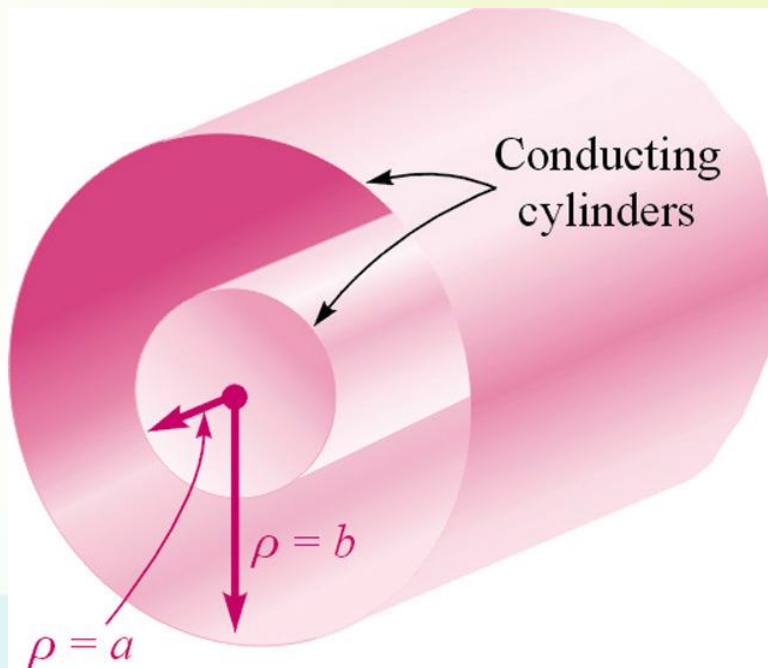
Gauss's law simplifies the task of finding  $\mathbf{D}$  near an infinite line charge.



$$\begin{aligned} Q &= \oint_{\text{cyl}} \mathbf{D}_S \cdot d\mathbf{S} = D_S \int_{\text{sides}} dS + 0 \int_{\text{top}} dS + 0 \int_{\text{bottom}} dS \\ &= D_S \int_{z=0}^L \int_{\phi=0}^{2\pi} \rho d\phi dz = D_S 2\pi\rho L \\ D_S = D_\rho &= \frac{Q}{2\pi\rho L} = \frac{\rho_L}{2\pi\rho} \end{aligned}$$



Infinite coaxial cable:

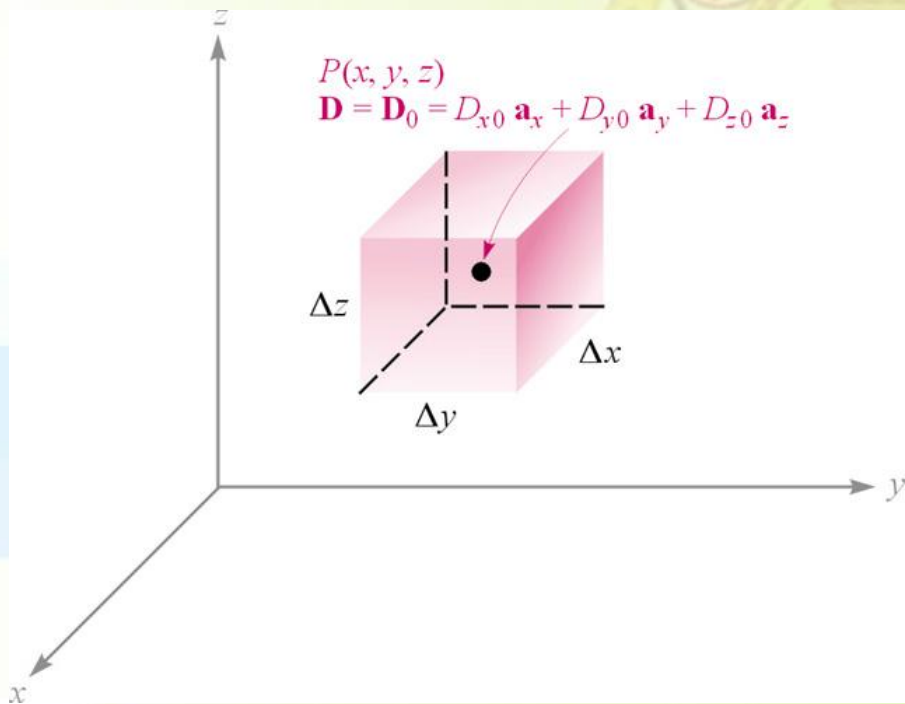


$$D = \frac{\rho L}{2\pi\rho} \mathbf{a}_\rho \quad (a < \rho < b)$$

$$D = 0 \quad (\rho > b)$$

# Differential Volume Element

- If we take a small enough closed surface, then  $\mathbf{D}$  is almost constant over the surface.



$$\oint_S \mathbf{D} \cdot d\mathbf{S} = \int_{\text{front}} + \int_{\text{back}} + \int_{\text{left}} + \int_{\text{right}} + \int_{\text{top}} + \int_{\text{bottom}}$$

$$\int_{\text{front}} \doteq \left( D_{x0} + \frac{\Delta x}{2} \frac{\partial D_x}{\partial x} \right) \Delta y \Delta z$$

$$\int_{\text{back}} \doteq \left( -D_{x0} + \frac{\Delta x}{2} \frac{\partial D_x}{\partial x} \right) \Delta y \Delta z$$

$$\int_{\text{front}} + \int_{\text{back}} \doteq \frac{\partial D_x}{\partial x} \Delta x \Delta y \Delta z$$

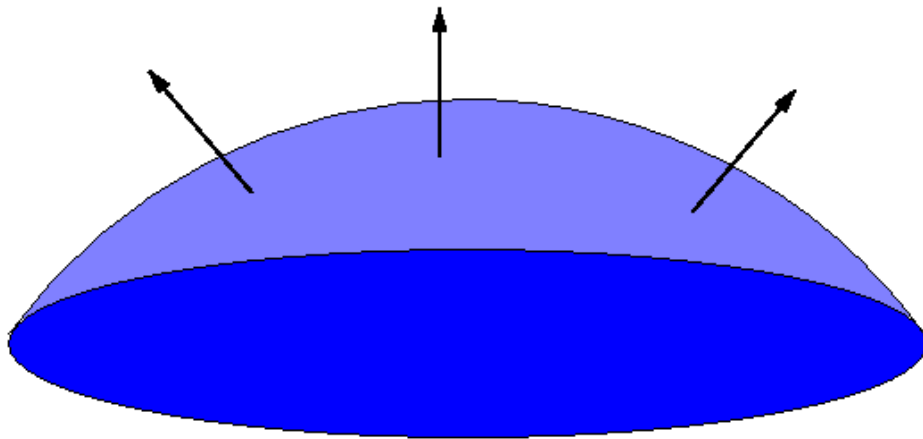
⋮

$$\oint_S \mathbf{D} \cdot d\mathbf{S} \doteq \left( \frac{\partial D_x}{\partial x} + \frac{\partial D_y}{\partial y} + \frac{\partial D_z}{\partial z} \right) \Delta x \Delta y \Delta z$$

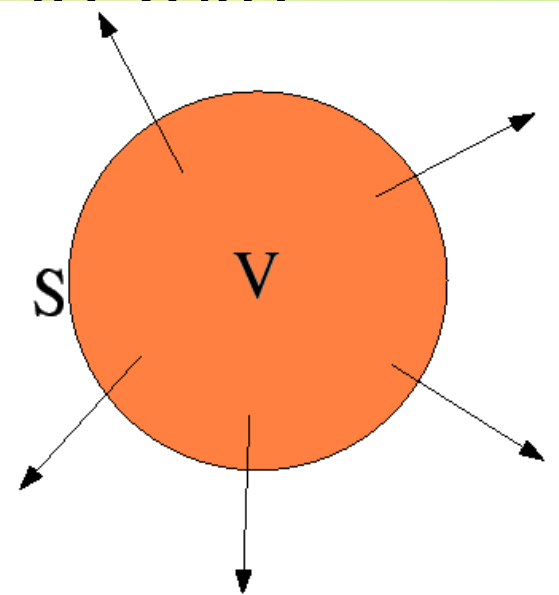
$$\text{Charge enclosed in volume } \Delta v \doteq \left( \frac{\partial D_x}{\partial x} + \frac{\partial D_y}{\partial y} + \frac{\partial D_z}{\partial z} \right) \times \text{volume } \Delta v$$

# Divergence

Divergence is the outflow of flux from a small closed surface area (per unit volume) as volume shrinks to zero



open surface.



closed surface

## -Water leaving a bathtub

- Closed surface (water itself) is essentially incompressible
- Net outflow is zero

## -Air leaving a punctured tire

- Divergence is positive, as closed surface (tire) exhibits net outflow



# Mathematical definition of divergence

$$\text{div}(\mathbf{D}) = \lim_{\Delta v \rightarrow 0} \int \frac{\mathbf{D}}{\Delta v} d\mathbf{S}$$

Surface integral as the volume element ( $\Delta v$ ) approaches zero

$\mathbf{D}$  is the vector flux density

$$\text{div}(\mathbf{D}) = \left( \frac{\delta D_x}{\delta x} + \frac{\delta D_y}{\delta y} + \frac{\delta D_z}{\delta z} \right)$$

- Cartesian

# Divergence in Other Coordinate Systems

## *Cylindrical*

$$\operatorname{div}(\mathbf{D}) = \frac{1}{\rho} \cdot \frac{\delta}{\delta \rho} (\rho \cdot \mathbf{D}_\rho) + \frac{1}{\rho} \cdot \frac{\delta \mathbf{D}_\phi}{\delta \phi} + \frac{\delta \mathbf{D}_z}{\delta z}$$

## *Spherical*

$$\operatorname{div}(\mathbf{D}) = \frac{1}{r^2} \cdot \frac{\delta (\mathbf{D}_r \cdot r^2)}{\delta r} + \frac{1}{r \cdot \sin(\theta)} \cdot \frac{\delta (\mathbf{D}_\theta \cdot \sin(\theta))}{\delta \theta} + \frac{1}{r \cdot \sin(\theta)} \cdot \frac{\delta \mathbf{D}_\phi}{\delta \phi}$$

# 4.1 Energy to move a point charge through a Field

- Force on  $Q$  due to an electric field

$$\underline{F_E = QE}$$

- Differential work done by an external source moving  $Q$

$$\underline{dW = -QE \cdot dL}$$

- Work required to move a charge a finite distance

$$W = -Q \int_{\text{init}}^{\text{final}} \mathbf{E} \cdot d\mathbf{L}$$



# Line Integral

- Work expression without using vectors

EL is the component of E in the dL direction

$$W = -Q \cdot \int_{\text{initial}}^{\text{final}} E_L dL$$

$$d\mathbf{L} = dx \mathbf{a}_x + dy \mathbf{a}_y + dz \mathbf{a}_z \quad (\text{cartesian})$$

$$d\mathbf{L} = d\rho \mathbf{a}_\rho + \rho d\phi \mathbf{a}_\phi + dz \mathbf{a}_z \quad (\text{cylindrical})$$

$$d\mathbf{L} = dr \mathbf{a}_r + r d\theta \mathbf{a}_\theta + r \sin \theta d\phi \mathbf{a}_\phi \quad (\text{spherical})$$

- Uniform electric field density

$$W = -QE \cdot L_{BA}$$

# Potential

- Measure potential difference between a point and something which has zero potential “ground”

$$V_{AB} = V_A - V_B$$

# Potential Field of a Point Charge

- Let  $V=0$  at infinity

$$V = \frac{Q}{4\pi\epsilon_0 r}$$

- Equipotential surface:
  - A surface composed of all points having the same potential

# Potential due to $n$ point charges

Continue adding charges

$$V(r) = \frac{Q_1}{4 \cdot \pi \cdot \epsilon_0 \cdot |r - r_1|} + \frac{Q_2}{4 \cdot \pi \cdot \epsilon_0 \cdot |r - r_2|} + \dots + \frac{Q_n}{4 \cdot \pi \cdot \epsilon_0 \cdot |r - r_n|}$$

$$V(r) = \sum_{m=1}^n \frac{Q_m}{4 \cdot \pi \cdot \epsilon_0 \cdot |r - r_m|}$$

# Potential as point charges become infinite

Volume of charge

$$V(r) = \int \frac{\rho_v(r_{\text{prime}})}{4 \cdot \pi \cdot \epsilon_0 \cdot |r - r_{\text{prime}}|} dv_{\text{prime}}$$

Line of charge

$$V(r) = \int \frac{\rho_L(r_{\text{prime}})}{4 \cdot \pi \cdot \epsilon_0 \cdot |r - r_{\text{prime}}|} dL_{\text{prime}}$$

Surface of charge

$$V(r) = \int \frac{\rho_S(r_{\text{prime}})}{4 \cdot \pi \cdot \epsilon_0 \cdot |r - r_{\text{prime}}|} dS_{\text{prime}}$$

# Potential as point charges become infinite

Volume of charge

$$V(r) = \int \frac{\rho_v(r_{\text{prime}})}{4 \cdot \pi \cdot \epsilon_0 \cdot |r - r_{\text{prime}}|} dv_{\text{prime}}$$

Line of charge

$$V(r) = \int \frac{\rho_L(r_{\text{prime}})}{4 \cdot \pi \cdot \epsilon_0 \cdot |r - r_{\text{prime}}|} dL_{\text{prime}}$$

Surface of charge

$$V(r) = \int \frac{\rho_S(r_{\text{prime}})}{4 \cdot \pi \cdot \epsilon_0 \cdot |r - r_{\text{prime}}|} dS_{\text{prime}}$$

# Gradients in different coordinate systems

The following equations are found on page 104 and inside the back cover of the text:

$$\text{grad}V = \frac{\delta V}{\delta x} \cdot a_x + \frac{\delta V}{\delta y} \cdot a_y + \frac{\delta V}{\delta z} \cdot a_z$$

Cartesian

$$\text{grad}V = \frac{\delta V}{\delta \rho} \cdot a_\rho + \frac{1}{\rho} \cdot \frac{\delta V}{\delta \phi} \cdot a_\phi + \frac{\delta V}{\delta z} \cdot a_z$$

Cylindrical

$$\text{grad}V = \frac{\delta V}{\delta r} \cdot a_r + \frac{1}{r} \cdot \frac{\delta V}{\delta \theta} \cdot a_\theta + \frac{1}{r \cdot \sin(\theta)} \cdot \frac{\delta V}{\delta \phi} \cdot a_\phi$$

Spherical

## Chapter 7 – Poisson's and Laplace Equations

A useful approach to the calculation of electric potentials  
Relates potential to the charge density.

The electric field is related to the charge density by the divergence relationship

$$\nabla \cdot E = \frac{\rho}{\epsilon_0}$$

$E$  = electric field  
 $\rho$  = charge density  
 $\epsilon_0$  = permittivity

The electric field is related to the electric potential by a gradient relationship

$$E = -\nabla V$$

Therefore the potential is related to the charge density by Poisson's equation

$$\nabla \cdot \nabla V = \nabla^2 V = \frac{-\rho}{\epsilon_0}$$

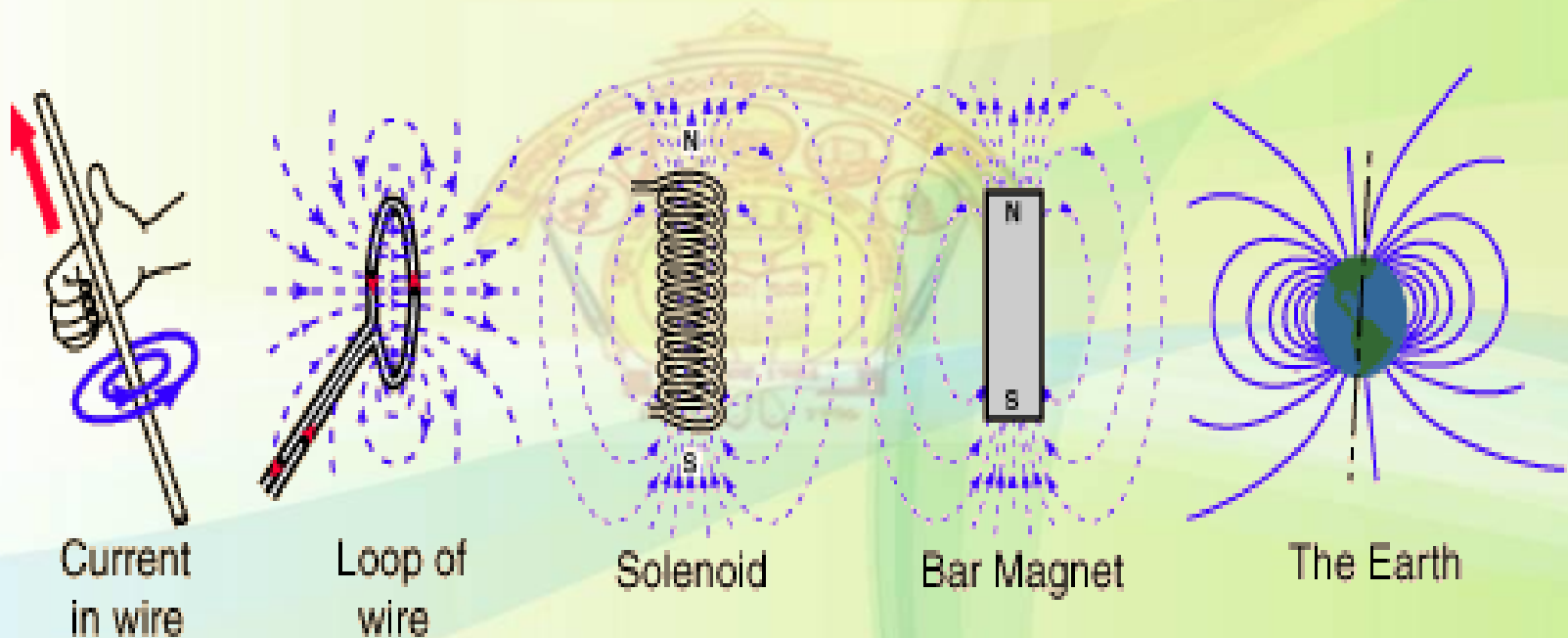
In a charge-free region of space, this becomes Laplace's equation

$$\nabla^2 V = 0$$



# Magnetic Field Sources

Magnetic fields are produced by electric currents, which can be macroscopic currents in wires, or microscopic currents associated with electrons in atomic orbits.



Magnetic Field Sources

# Maxwell's equations

$$\nabla \times \underline{E} = -\frac{\partial \underline{B}}{\partial t} - \underline{K}_c - \underline{K}_i$$

$$\nabla \times \underline{H} = \frac{\partial \underline{D}}{\partial t} + \underline{J}_c + \underline{J}_i$$

$$\nabla \cdot \underline{D} = q_{ev}$$

$$\nabla \cdot \underline{B} = q_{mv}$$

# Maxwell's equations for TVF

| Differential form                                 | Controlling principle                    | Integral form  |       |
|---|--|--|-------|
| $\nabla \times \vec{H} = \dot{\vec{D}} + \vec{J}$ | Ampere's Circuital Law                   | $\oint \vec{H} \cdot d\vec{L} = \int \dot{\vec{D}} + \vec{J} \cdot d\vec{S}$ | (I)   |
| $\nabla \times \vec{E} = -\dot{\vec{B}}$          | Potential around a closed path is zero   | $\oint \vec{E} \cdot d\vec{L} = -\int \dot{\vec{B}} \cdot d\vec{S}$          | (II)  |
| $\nabla \cdot \vec{D} = \rho$                     | Gauss's Law                              | $\oint \vec{D} \cdot d\vec{S} = \int \rho \, dv$                             | (III) |
| $\nabla \cdot \vec{B}$                            | Non-existence of isolated magnetic poles | $\oint \vec{B} \cdot d\vec{S} = 0$   | (IV)  |

# Queries ....?

