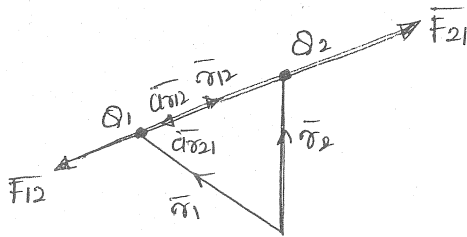


Fourth Semester BE. Degree Examination, June/July 2013
Field Theory

- (1) a) State and explain Coulomb's Law of force between two point charges also indicate the units of quantities in the force equation

→ It states that "The electrostatic force of attraction or repulsion between the two point charges is directly proportional to the product of magnitude of charges and is inversely proportional to the square of distance between them."



Consider the two point charges Q_1 and Q_2 as shown in fig separated by distance 'r'. The charge Q_1 exerts a force on Q_2 while Q_2 also exerts force on Q_1 the force acting along line joining Q_1 and Q_2 the force exerted betⁿ them is repulsive if the polarities are of same type and attractive if charge polarities are different

$$F \propto Q_1 Q_2$$

$$F \propto \frac{1}{|r_{12}|^2}$$

$$F \propto \frac{Q_1 Q_2}{|r_{12}|^2}$$

$$F = K \frac{Q_1 Q_2}{|r_{12}|^2}$$

where K is constant of proportionality and value is $\frac{1}{4\pi\epsilon}$

$$\therefore F = \frac{Q_1 Q_2}{4\pi\epsilon |r_{12}|^2}$$

Inverted form $F = \frac{Q_1 Q_2}{4\pi\epsilon |r_{12}|^2} \hat{a}_{r_{12}}$

From above fig we have

$$\bar{r}_1 + \bar{r}_{12} = \bar{r}_2$$

$$\therefore \bar{r}_{12} = \bar{r}_2 - \bar{r}_1$$

$$\text{and, } |\bar{r}_{12}| = \sqrt{r_2^2 + r_1^2}$$

$$\text{Now } \frac{\bar{r}_{12}}{|\bar{r}_{12}|} = \frac{\bar{r}_2 - \bar{r}_1}{\sqrt{r_2^2 + r_1^2}}$$

$$\text{Then } \vec{F} = \frac{Q_1 Q_2}{4\pi\epsilon (\sqrt{r_2^2 + r_1^2})^2} \cdot \frac{\bar{r}_2 - \bar{r}_1}{\sqrt{r_2^2 + r_1^2}}$$

$$\therefore \boxed{\vec{F}_{21} = \frac{Q_1 Q_2}{4\pi\epsilon (r_2^2 + r_1^2)^{3/2}} (\bar{r}_2 - \bar{r}_1)} \quad \text{--- (1)}$$

|||^{ly} Force on Q_1 due to Q_2 is given by

$$\boxed{\vec{F}_{12} = \frac{Q_1 Q_2}{4\pi\epsilon (r_1^2 + r_2^2)^{3/2}} (\bar{r}_1 - \bar{r}_2)} \quad \text{--- (2)}$$

$$\text{From (1) \& (2) } \vec{F}_{21} = -\vec{F}_{12}$$

Units :

$$Q_1, Q_2 = \text{coulombs}$$

$$|r_1| = \text{meters}$$

$$\epsilon_0 = \text{Absolute Permittivity} = 8.85 \times 10^{-12} \text{ F/m}$$

$$\epsilon_r = \text{Relative Permittivity value depends on type of medium}$$

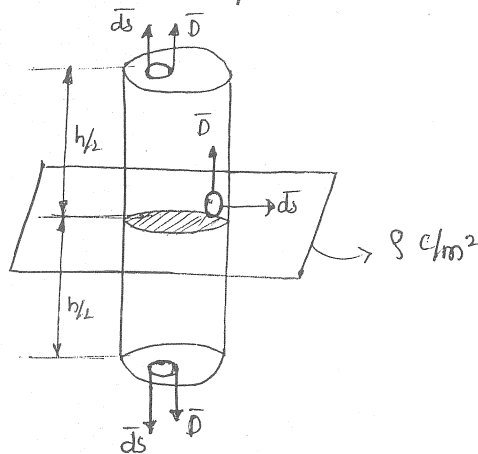
$$\epsilon_r = 1 \text{ (for air medium).}$$

b) State and apply Gauss law to obtain the expression for electric field intensity due to an infinite sheet charge with a surface charge density $\rho \text{ C/m}^2$ and area $A. \text{m}^2$

→ Gauss Law: "The total normal electric flux passing through any closed surface is equal to the charge enclosed by that surface"

$$\oint_S \vec{D} \cdot d\vec{s} = Q_{\text{enclosed}}$$

Electric Field Intensity due to infinite sheet charge:



Consider an uniform sheet charge with surface charge density of $\rho \text{ C/m}^2$ and area as $A \cdot \text{m}^2$

∴ From Gauss Law $\oint_S \vec{D} \cdot d\vec{s} = Q_{\text{enclosed}}$

For different surfaces $\int_{TS} \vec{D} \cdot d\vec{s} + \int_{BS} \vec{D} \cdot d\vec{s} + \int_{CS} \vec{D} \cdot d\vec{s} = Q_{\text{enclosed}}$ with cylindrical co-s

$$\int_{TS} D_z \vec{a}_z \cdot r dr d\phi \vec{a}_z + \int_{BS} D_z (-\vec{a}_z) \cdot r dr d\phi (-\vec{a}_z) + 0 = Q_{\text{enclosed}}$$

$$\int_S D_z r dr d\phi + \int_S D_z r dr d\phi = Q_{\text{enclosed}}$$

$$D_z \int_{r=0}^A r dr \int_0^{2\pi} d\phi + D_z \int_{r=0}^A r dr \int_0^{2\pi} d\phi = Q_{\text{enclosed}}$$

$$D_z \left[\frac{r^2}{2} \right]_0^A \left[\phi \right]_0^{2\pi} + D_z \left[\frac{r^2}{2} \right]_0^A \left[\phi \right]_0^{2\pi} = Q_{\text{enclosed}}$$

$$D_z \left(\frac{A^2}{2} \right) 2\pi + D_z \left(\frac{A^2}{2} \right) 2\pi = Q_{\text{enclosed}}$$

$$D_z 2\pi A^2 = \rho \pi A^2$$

$$D_z = \frac{\rho}{2}$$

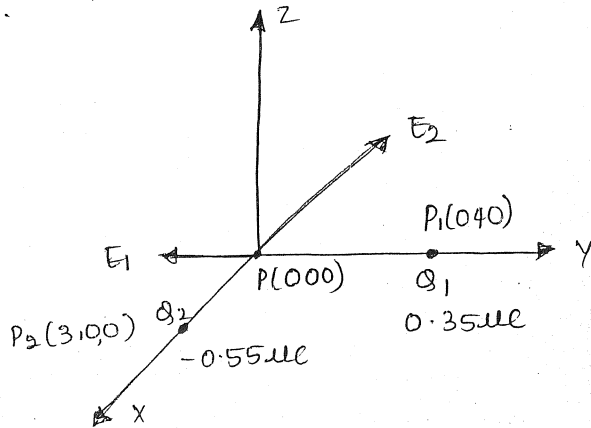
$$\bar{D} = D_2 \bar{a}_2$$

$$\therefore \bar{D} = \frac{\rho}{2} \bar{a}_2$$

WKT $\bar{D} = \epsilon \bar{E}$

$$\therefore \bar{E} = \frac{\rho}{2\epsilon} \bar{a}_2$$

c) Find: i) Electric Field intensity and ii) Electric Flux density at the origin due to $Q_1 = 0.35 \mu\text{C}$ at $(0, 4, 0)\text{m}$ and $Q_2 = -0.55 \mu\text{C}$ at $(3, 0, 0)\text{m}$.



$$\bar{E} = \frac{Q}{4\pi\epsilon r^2} \bar{a}_r \quad \text{From electric field intensity}$$

$$\bar{E}_1 = \frac{Q_1}{4\pi\epsilon |\bar{P}_1P|^2} \bar{a}_{P_1P}$$

$$\bar{P}_1P = 0\bar{a}_x - 4\bar{a}_y + 0\bar{a}_z = -4\bar{a}_y$$

$$|\bar{P}_1P| = 4$$

$$\bar{E}_1 = \frac{0.35 \times 10^{-6}}{4\pi \times 8.85 \times 10^{-12} \times 4^2} (-\bar{a}_y)$$

$$\therefore \bar{a}_{P_1P} = \frac{\bar{P}_1P}{|\bar{P}_1P|} = \frac{-4\bar{a}_y}{4} = -\bar{a}_y$$

$$\bar{E}_1 = -196.6 \bar{a}_y \text{ N/C}$$

$$\bar{E}_2 = \frac{Q_2}{4\pi\epsilon |\bar{P}_2P|^2} \bar{a}_{P_2P}$$

$$\bar{P}_2P = -3\bar{a}_x + 0\bar{a}_y + 0\bar{a}_z = -3\bar{a}_x$$

$$|\bar{P}_2P| = 3$$

$$\bar{E}_2 = \frac{-0.55 \times 10^{-6}}{4\pi \times 8.85 \times 10^{-12} \times 3^2} (-\bar{a}_x)$$

$$\bar{a}_{P_2P} = \frac{\bar{P}_2P}{|\bar{P}_2P|} = \frac{-3\bar{a}_x}{3} = -\bar{a}_x$$

$$\bar{E}_2 = 549.42 \bar{a}_x \text{ N/C}$$

We have the relation $\vec{D} = \epsilon \vec{E}$ or $\vec{D} = \frac{Q}{4\pi r^2} \vec{a}_r$

Magnetic Flux density $\vec{D} = \epsilon \vec{E}$

$$\therefore \vec{D}_1 = \epsilon \vec{E}_1$$

$$= 8.85 \times 10^{12} \times (-196.6 \vec{a}_y)$$

$$\vec{D}_1 = -1.744 \times 10^{-9} \vec{a}_y$$

$$D_2 = \epsilon E_2$$

$$= 8.85 \times 10^{12} \times 549.42 \vec{a}_x$$

$$= 4.86 \times 10^{-9} \vec{a}_x$$

$$\therefore \vec{D} = \vec{D}_1 + \vec{D}_2$$

$$\vec{D} = (4.86 \vec{a}_x - 1.744 \vec{a}_y) \times 10^{-9}$$

$$|\vec{D}| = \sqrt{4.86^2 + (-1.744)^2} \times (10^{-9})^2$$

$$|\vec{D}| = ~~8.22~~ 5.1620 \times 10^{-9} \text{ C/m}^2$$

- (2) a) Explain with mathematical expressions i) Potential difference
ii) Absolute potential. iii) Potential gradient.

→ i) Potential Difference

In electric field, work done in moving a charge from B to A is given by

$$W = -Q \int_B^A \vec{E} \cdot d\vec{l}$$

If charge is selected as unit test charge then from above eqⁿ we get the work done in moving a unit charge from B to A in the field \vec{E} this work done in moving a unit charge from B to A in the field \vec{E} is called Potential Difference between the point B and A is denoted as V

$$\therefore \text{Potential Difference } V = - \int_B^A \vec{E} \cdot d\vec{l}$$

Thus "work done per unit charge in moving unit charge from B to A in the field \vec{E} is called potential difference between the point B & A

$$V_{AB} = - \int_B^A \vec{E} \cdot d\vec{l}$$

ii) Absolute potential

Absolute potential can be expressed at various points in the field such absolute potential are measured wrt specified reference position, such a reference position is assumed to be at zero

Consider a potential difference due to movement of unit charge B to A in a field of point charge is given by

$$V_{AB} = \frac{Q}{4\pi\epsilon_0} \left(\frac{1}{r_A} - \frac{1}{r_B} \right)$$

Now if the charge is moved from point ∞ to A, i.e. $r_B = \infty$

$$\frac{1}{r_B} = \frac{1}{\infty} = 0$$

$$\therefore V_{AB} = \frac{Q}{4\pi\epsilon_0 r_A} \text{ v} \quad \text{--- potential at point A}$$

which is also called absolute potential of point A

iii) absolute potential of point B is given as

$$V_B = \frac{Q}{4\pi\epsilon_0 r_B} \text{ v}$$

iii) Potential Gradient

Consider an electric field due to a positive charge placed at the origin of sphere then

$$V = - \int \vec{E} \cdot d\vec{l} = \frac{Q}{4\pi\epsilon_0 r}$$

The potential decreases as distance of point from the charge increases $\therefore V_{AB} = \Delta V = - \vec{E} \cdot d\vec{l}$

The rate of change of potential wrt the distance is called the potential gradient

$$\frac{dV}{dL} = \lim_{\Delta L \rightarrow 0} \frac{\Delta V}{\Delta L} = \text{Potential Gradient}$$

(b) Derive an expression for equation of continuity of current

→ The continuity equation of current is based on the principle of conservation of charge which states that "charge can neither be created nor be destroyed"

Consider a closed surface 'S' with ρ density \vec{J} then the total ρ crossing the surface is given as

$$I = \oint_S \vec{J} \cdot d\vec{s} \quad \text{--- (1)}$$

It has been mentioned earlier that the ρ means the flow of +ve charges hence I is constituted due to outward flow of +ve charges from closed surface 'S' according to the principle of conservation of

Q_i = charge within closed surface

$-\frac{dQ_i}{dt}$ = rate of decrease of charge inside the closed surface

The -ve charge indicates decrease in charge

Due to principle of conservation of charge this rate of decrease is same as rate of outward flow of charge which is a ρ

$$I = \oint_S \vec{J} \cdot d\vec{s} = -\frac{dQ_i}{dt} \quad \text{(2) Integral form of } \rho \text{ continuity}$$

The negative sign in eqⁿ indicates outward flow of ρ from closed surface so the eqⁿ (2) indicating outward flow of ρ

$$\oint_S \vec{J} \cdot d\vec{s} = -I = +\frac{dQ_i}{dt}$$

The point form of continuity equation can be obtained from integral form by using divergence theorem

$$\oint_S \vec{J} \cdot d\vec{s} = \int_{\text{vol}} (\nabla \cdot \vec{J}) dv$$

$$-\frac{dQ_i}{dt} = \int_V (\nabla \cdot \vec{J}) dv$$

But $Q_i = \int_V \rho_v dv$

where ρ_v = volume charge density

$$\therefore \int_V (\nabla \cdot \vec{J}) dv = \frac{d}{dt} \left[\int_V \rho_v dv \right] = - \int_V \frac{\partial \rho_v}{\partial t} dv$$

for constant surface the derivative becomes partial derivative

$$\therefore \int_V (\nabla \cdot \vec{J}) dv = \int_V - \frac{\partial \rho_v}{\partial t} dv$$

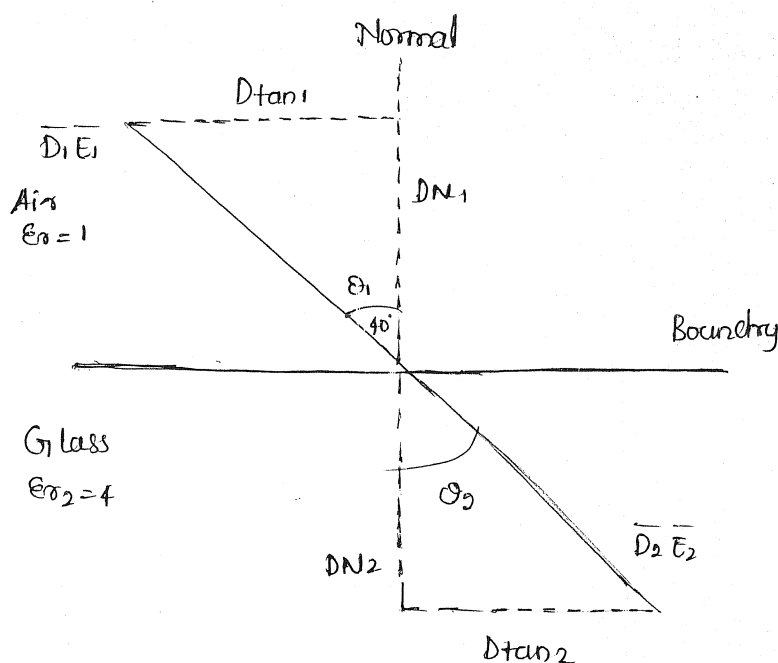
If the relation for any volume is true for incremental volume

$$(\nabla \cdot \vec{J}) dv = \frac{\partial \rho_v}{\partial t} dv$$

$$\boxed{\nabla \cdot \vec{J} = \frac{\partial \rho_v}{\partial t}}$$

This is point form of continuity equation of current

- c) At the boundary between glass ($\epsilon_r = 4$) and air the lines of electric field make an angle of 40° with normal to the boundary. If the electric flux density in air is $0.25 \mu\text{C}/\text{m}^2$, determine the orientation of magnitude of i) electric flux density and ii) electric field intensity in glass



The arrangement as shown in above fig For boundary betw
the two dielectrics

$$DN_1 = DN_2$$

$$\text{and } \frac{EN_1}{EN_2} = \frac{\epsilon_2}{\epsilon_1}$$

$$\text{Now } D_1 = 0.25 \mu\text{C}/\text{m}^2$$

$$\cos \theta_1 = \frac{DN_1}{D_1}$$

$$\cos 40^\circ = \frac{DN_1}{0.25 \mu}$$

$$\therefore DN_1 = 0.1915 \mu\text{C}/\text{m}^2$$

$$\text{As } DN_2 = DN_1$$

$$\therefore DN_2 = 0.1915 \mu\text{C}/\text{m}^2$$

$$D_1 = \sqrt{(DN_1)^2 + (D \tan \theta_1)^2}$$

$$0.25 = \sqrt{(0.1915)^2 + (D \tan \theta_1)^2}$$

$$\therefore D \tan \theta_1 = 0.1607 \mu\text{C}/\text{m}^2$$

$$\text{But } \frac{D \tan \theta_1}{D \tan \theta_2} = \frac{\epsilon_1}{\epsilon_2} = \frac{\epsilon_{r1}}{\epsilon_{r2}}$$

$$D \tan \theta_2 = \frac{\epsilon_{r2}}{\epsilon_{r1}} \times D \tan \theta_1 = \frac{4}{1} \times 0.1607 = 0.6428 \mu\text{C}/\text{m}^2$$

$$D_2 = \sqrt{(DN_2)^2 + (D \tan \theta_2)^2}$$

$$D_2 = 0.6707 \mu\text{C}/\text{m}^2$$

$$\text{and } \cos \theta_2 = \frac{DN_2}{D_2} = \frac{0.1915}{0.6707} = 0.2855$$

$$\therefore \theta_2 = 73.41^\circ$$

3 (a) Derive Poisson's and Laplace equations starting from Gauss law

→ When charge and potentials are given at same boundary regions then potential and electric field intensity at any point may be obtained by using Poisson's and Laplace equations

Consider Gauss Law in point form

$$\nabla \cdot \mathbf{D} = \rho_v \quad \text{--- (1)}$$

where ρ_v = volume charge density and

\mathbf{D} = Electric flux density

WKT Homogeneous Isotropic equation $|\mathbf{D}| = \epsilon \mathbf{E}$

then eqn (1) becomes

$$\nabla \cdot \epsilon \mathbf{E} = \rho_v$$

$$\nabla \cdot \mathbf{E} = \frac{\rho_v}{\epsilon}$$

we have $\mathbf{E} = -\nabla V$

$$\therefore \nabla \cdot (-\nabla V) = \frac{\rho_v}{\epsilon}$$

$$-\nabla^2 V = \frac{\rho_v}{\epsilon}$$

$$\boxed{\nabla^2 V = -\frac{\rho_v}{\epsilon}} \quad \text{--- (2)}$$

Equation (2) is called Poisson's equation.

For dielectric medium $\rho_v = 0$ then eqn (2) can be rewritten as

$$\boxed{\nabla^2 V = 0} \quad \text{--- (3)}$$

equation (3) is called Laplace equation

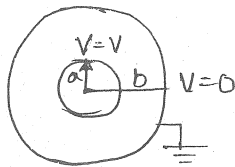
where ∇^2 = Laplace operator

$$\nabla^2 v = \frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2} + \frac{\partial^2 v}{\partial z^2} \longrightarrow \text{Rectangular coordinate}$$

$$\nabla^2 v = \frac{1}{s} \frac{\partial}{\partial s} \left(s \frac{\partial v}{\partial s} \right) + \frac{1}{s} \frac{\partial^2 v}{\partial \phi^2} + \frac{\partial^2 v}{\partial z^2} \longrightarrow \text{Cylindrical coordinate}$$

$$\nabla^2 v = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial v}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial v}{\partial \theta} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial^2 v}{\partial \phi^2} \longrightarrow \text{Spherical coordinate}$$

- (b) Using Laplace equation derive an expression for capacitance of concentric spherical capacitor the inner spherical conductor is of radius 'a' and potential 'V' while outer conductor of radius 'b' and potential zero.



The concentric conductor are shown above fig at $r=b$ and $v=0$ hence outer potential is shown at zero potential

The electric field intensity E will be only in only radial direction hence v is changing only in radial distance r and not the function of θ & ϕ

equation for $\nabla^2 v$ is spherical coordinate.

$$\nabla^2 v = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial v}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial v}{\partial \theta} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial^2 v}{\partial \phi^2} = 0$$

As there is no potential variation along θ and ϕ axis \therefore

$$\frac{\partial v}{\partial \theta} = \frac{\partial v}{\partial \phi} = 0$$

From Laplace equation $\nabla^2 v = 0$

$$\text{i.e. } \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial v}{\partial r} \right) = 0$$

$$\frac{\partial}{\partial r} \left(r^2 \frac{\partial v}{\partial r} \right) = 0$$

Integrating on both sides wrt r

$$\int \frac{\partial}{\partial r} \left(r^2 \frac{\partial V}{\partial r} \right) = 0$$

$$\therefore r^2 \frac{\partial V}{\partial r} = C_1$$

$$\frac{\partial V}{\partial r} = \frac{C_1}{r^2}$$

Integrating again wrt r on b.s

$$V = -\frac{C_1}{r} + C_2 \quad \text{--- (1)}$$

Substituting the boundary conditions

$$r = a, V = V$$

$$V = -\frac{C_1}{a} + C_2$$

$$r = b, V = 0$$

$$0 = -\frac{C_1}{b} + C_2$$

$$V = -C_1 \left(\frac{1}{a} - \frac{1}{b} \right)$$

$$-V = -C_1 \left(\frac{1}{b} - \frac{1}{a} \right)$$

$$\therefore C_1 = \frac{V(ab)}{(a-b)}$$

$$C_2 = \frac{C_1}{b} = \frac{V(ab)}{b(a-b)}$$

Substituting in potential equation (1)

$$V = \frac{-V(ab)}{(a-b)r} + \frac{V(ab)}{b(a-b)}$$

we have $\vec{E} = -\nabla V$

$$\therefore \vec{E} = - \left[\frac{\partial V}{\partial r} \vec{a}_r + \frac{1}{r} \frac{\partial V}{\partial \theta} \vec{a}_\theta + \frac{1}{r \sin \theta} \frac{\partial V}{\partial \phi} \vec{a}_\phi \right]$$

as field is varying only along r -axis $\therefore \theta = \phi = 0$

$$\therefore \vec{E} = - \left[\frac{\partial}{\partial r} \left[\frac{-V(ab)}{r(a-b)} + \frac{V(a)}{(a-b)} \right] \vec{a}_r \right]$$

$$\therefore \vec{E} = - \left[\frac{V(ab)}{(a-b)r^2} + 0 \right] \vec{ar}$$

$$\vec{E} = \frac{-V(ab)}{r^2(a-b)} \vec{ar}$$

$$|\vec{E}| = \frac{V(ab)}{r^2(a-b)} \vec{ar}$$

$$\text{WKT } Q = \frac{Q_s}{\epsilon}$$

$$\text{also } Q = \int \vec{D} \cdot \vec{A}$$

$$Q = \epsilon \vec{E} \cdot \vec{A}$$

$$\therefore Q = \epsilon \frac{V(ab)}{r^2(a-b)} 4\pi r^2$$

$$Q = \frac{4\pi\epsilon V(ab)}{(a-b)}$$

$$\text{capacitance } C = \frac{Q}{V_0} = \frac{4\pi\epsilon V(ab)}{V_0(a-b)}$$

(c) Determine whether or not following potential fields satisfy Laplace equation (i) $v = 2x^2 - 3y^2 + z^2$ (ii) $r^2 + z^2$

$$\rightarrow v = 2x^2 - 3y^2 + z^2$$

$$\nabla^2 v = \frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2} + \frac{\partial^2 v}{\partial z^2} \rightarrow \text{Rectangular coordinate}$$

$$\nabla^2 v = \frac{\partial^2}{\partial x^2} (2x^2 - 3y^2 + z^2) + \frac{\partial^2}{\partial y^2} (2x^2 - 3y^2 + z^2) + \frac{\partial^2}{\partial z^2} (2x^2 - 3y^2 + z^2)$$

$$= \frac{\partial}{\partial x} (4x) + \frac{\partial}{\partial y} (-6y) + \frac{\partial}{\partial z} (2z)$$

$$= 4 - 6 + 2$$

$$\nabla^2 v = 0$$

From Laplace equation $\nabla^2 v = 0$

It satisfies Laplace equation.

$$v = r^2 + z^2$$

$$\nabla^2 v = \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial v}{\partial r} \right) + \frac{1}{r} \left(\frac{\partial^2 v}{\partial \phi^2} \right) + \frac{\partial^2 v}{\partial z^2} \longrightarrow \text{Cylindrical coordinates}$$

$$= \frac{1}{r} \frac{\partial}{\partial r} (r \times 2r) + \frac{1}{r} (0) + 2$$

$$= \frac{1}{r} \times 4r + 0 + 2$$

$$\nabla^2 v = 6$$

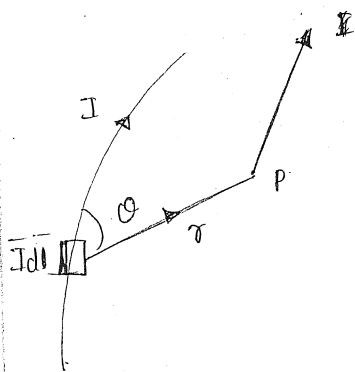
As Laplace equation $\nabla^2 v = 0$

So it does not satisfy Laplace equation

4(a) Write Explanatory note on Biot Savart Law

→ This Law states that "The magnetic field intensity at point produced by differential ϕ/n element is

- directly proportional to the product of ϕ/n I & differential ϕ/n element $d\vec{l}$
- is directly proportional to the sin of angle between the element and the line joining point 'p' to element and
- is Inversely proportional to the square of distance R betⁿ the point P and ϕ/n element (i)



$$dH \propto I |d\vec{l}|$$

$$dH \propto \sin \theta$$

$$dH \propto \frac{1}{r^2}$$

$$dH \propto \frac{I |d\vec{l}| \sin \theta}{r^2}$$

$$dH = \frac{k I |d\vec{l}| \sin \theta}{r^2}$$

where k is constant of proportionality $k = \frac{1}{4\pi}$

$$\therefore dH = \frac{I |d\vec{l}| \sin \theta}{4\pi r^2} \quad \because \vec{a} = \frac{\vec{r}}{|\vec{r}|}$$

(b) Discuss the concept of scalar and vector magnetic potential and derive the expression for Poisson's equation in magnetostatics

→ Scalar Magnetic Potential

If V_m is the scalar magnetic potential then it must satisfy the equation $\nabla \times \nabla V_m = 0$ — (1)

But scalar magnetic potential is related to magnetic field intensity \vec{H} as $\vec{H} = -\nabla V_m$

using equation (1)

$$\nabla \times (-\vec{H}) = 0 \quad \text{i.e.} \quad \nabla \times \vec{H} = 0$$

$$\text{But} \quad \nabla \times \vec{H} = \vec{J} \quad \text{i.e.} \quad \vec{J} = 0$$

The scalar magnetic potential V_m can be defined as for source free region where \vec{J} is current density is zero

$$\vec{H} = -\nabla V_m \quad \text{only for } \vec{J} = 0$$

||^{br} to their relation between \vec{E} and electric scalar potential magnetic scalar potential can be expressed in terms of \vec{H} as

$$V_{m,ab} = - \int_b^a \vec{H} \cdot d\vec{l} \quad \text{--- specified path}$$

Vector Magnetic Potential

The magnetic potential denoted as \vec{A} and measured in Wb/m it has to satisfy the eqn that divergence of curl of a vector is always zero

$$\nabla \cdot (\nabla \times \vec{A}) = 0 \quad \text{--- } \vec{A} \text{ - Vector magnetic potential}$$

$$\nabla \cdot \vec{B} = 0 \quad \text{--- From Laplace in scalar magnetic potential}$$

$$\vec{B} = \nabla \times \vec{A}$$

Thus curl of vector magnetic potential is flux density

$$\nabla \times \vec{A} = \vec{J}$$

$$\nabla \times \frac{\mathbf{B}}{\mu_0} = \mathbf{J} \quad \text{---} \quad \left[\because \mathbf{B} = \mu_0 \mathbf{H} \right]$$

$$\nabla \times \mathbf{B} = \mu_0 \mathbf{J}$$

$$\nabla \times \nabla \times \mathbf{A} = \mu_0 \mathbf{J}$$

Using vector identity to express left hand side we can write

$$\nabla (\nabla \cdot \mathbf{A}) - \nabla^2 \mathbf{A} = \mu_0 \mathbf{J}$$

$$\therefore \mathbf{J} = \frac{1}{\mu_0} [\nabla \times \nabla \times \mathbf{A}] = \frac{1}{\mu_0} [\nabla (\nabla \cdot \mathbf{A}) - \nabla^2 \mathbf{A}] \quad \text{---} \quad (2)$$

Thus if vector magnetic potential is known then ∇ density \mathbf{J} can be obtained for defining \mathbf{A} the ∇ density need to zero

Poisson's Equation in Magnetostatics

In vector algebra the vector can fully defined if its curl and divergence are defined for a vector magnetic potential \mathbf{A} will be defined as $\nabla \times \mathbf{A} = \mathbf{B}$ which is known but to completely define \mathbf{A} its divergence must be known assume $\nabla \cdot \mathbf{A} = 0$ the divergence \mathbf{A} is zero this is consistent with some other conditions to be

Using eqⁿ (2)

$$\mathbf{J} = \frac{1}{\mu_0} [-\nabla^2 \mathbf{A}]$$

$$\boxed{\nabla^2 \mathbf{A} = -\mu_0 \mathbf{J}}$$

This is called Poisson's equation in Magnetostatics.

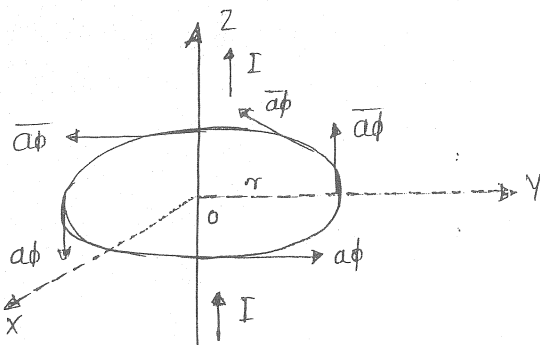
- (c) State and prove Ampere's circuit law and apply it to straight solid conductor to calculate the magnetic field intensity.

→ "The line integral of magnetic field intensity \mathbf{H} around a closed path is exactly equal to the direct c/n enclosed by that path" $\oint \mathbf{H} \cdot d\mathbf{l} = I_{\text{enclosed}}$

Proof

Consider a long straight conductor carrying $\phi_0 I$ placed along Z-axis as shown in below fig consider a circular closed path of radius r which encloses the straight conductor carrying $\phi_0 I$ the point P is a r^{th} distance from the conductor consider $d\vec{l}$ at point P which is in \vec{a}_ϕ direction tangential to circular path at P.

$$d\vec{l} = r d\phi \vec{a}_\phi$$



From Biot savart's Law

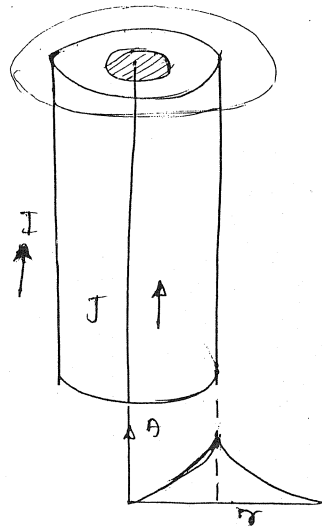
$$H = \frac{I}{2\pi r} \vec{a}_\phi$$

$$\begin{aligned} \oint \vec{H} \cdot d\vec{l} &= \int \frac{I}{2\pi r} \vec{a}_\phi (dr \vec{a}_r + r d\phi \vec{a}_\phi + dz \vec{a}_z) \\ &= \int \frac{I}{2\pi r} r d\phi \\ &= \frac{I}{2\pi} \int_0^{2\pi} d\phi \\ &= \frac{I}{2\pi} \times 2\pi \end{aligned}$$

$$\text{Thus } \boxed{\oint \vec{H} \cdot d\vec{l} = I \text{ enclosed}}$$

Magnetic field Intensity due to straight & solid conductor

Consider a straight solid conductor of length and it carrying the current of I with current density J . Let a be the cross-sectional area of solid conductor. Then magnetic field intensity can be obtained inside & outside the conductor as shown in below fig



$H = ?$ at $r < 0$ (H is inside the cond^r)

$$\oint \vec{H} \cdot d\vec{l} = I_{enc}$$

$$\oint H \phi \vec{a}_\phi \cdot r d\phi \vec{a}_\phi = I_{enc}$$

$$\oint H \phi \cdot r d\phi = I_{enc}$$

$$H \phi r \int_0^{2\pi} d\phi = I_{enc}$$

$$H \phi r 2\pi = J \times \pi r^2$$

$$H \phi = \frac{J r}{2\pi}$$

$$\vec{H} = H \phi \vec{a}_\phi$$

$$= \frac{J \cdot r}{2\pi} \vec{a}_\phi$$

$$\vec{H} = \frac{I}{2\pi a^2} \vec{a}_\phi$$

$H = ?$ at $r > 0$ (outside the cond^r)

$$\oint \vec{H} \cdot d\vec{l} = I_{enc}$$

$$\oint H \phi \vec{a}_\phi \cdot r d\phi \vec{a}_\phi = I_{enc}$$

$$\oint H \phi r d\phi = I_{enc}$$

$$H \phi r 2\pi = I_{enc}$$

$$H \phi = \frac{I}{2\pi r}$$

$$\vec{H} = H \phi \vec{a}_\phi$$

$$\vec{H} = \frac{I}{2\pi r} \cdot \vec{a}_\phi$$

(5) a) Find the force on expression for force on a differential γ_n carrying elements.

→ Force exerted on a differential γ_n element of charge dq moving in a steady magnetic field is given by

$$d\vec{F} = dq \vec{v} \times \vec{B} \quad \text{N} \quad \text{--- (1)}$$

The γ_n density \vec{J} can be expressed in terms of velocity of volume charge density as,

$$\vec{J} = \rho_v \vec{v} \quad \text{--- (2)}$$

But differential element charge can be expressed in terms of volume charge density as

$$dq = \rho_v dv$$

Substituting the value of dq in eqⁿ (1)

$$\therefore d\vec{F} = \rho_v dv \vec{v} \times \vec{B}$$

Expressing $d\vec{F}$ in terms of \vec{J} using eqⁿ (2) we can write.

$$d\vec{F} = \vec{J} \times \vec{B} \quad dv \quad \text{--- (3)}$$

we have the relation $\vec{J} dv = \vec{K} ds = I d\vec{l}$

Then force on differential γ_n density is given by

$$d\vec{F} = \vec{K} \times \vec{B} \quad ds \quad \text{--- (4)}$$

Similarly Force exerted on differential γ_n element is given by

$$d\vec{F} = (I d\vec{l} \times \vec{B})$$

Integrating eqⁿ (4) over volume.

$$\vec{F} = \int_V \vec{J} \times \vec{B} \quad dv$$

Integrating eqⁿ (5) over closed or open surface.

$$F = \int_S \vec{K} \times \vec{B} \quad ds$$

iii) integrating eqn (6) over closed path

$$\vec{F} = \oint I d\vec{l} \times \vec{B}$$

If the conductor is straight the \vec{B} is uniform along it then integrating eqn (6) we get simple expression for force as

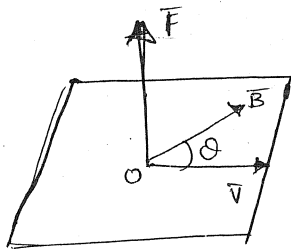
$$\boxed{\vec{F} = I \vec{L} \times \vec{B}}$$

magnitude of the Force is given by

$$\boxed{F = B I L \sin \theta}$$

(b) Define Lorentz force equation and mention the applications of its solution

→ Lorentz Force: It is the product of charge and sum of Electric field & cross product of velocity of charge with flux density (N)



We know that $\vec{F}_e = q\vec{E}$ — (1)

$$\vec{F}_m = q|\vec{v}||\vec{B}|\sin\theta$$

$$\vec{F}_m = q(\vec{v} \times \vec{B}) \text{ — (2)}$$

From (1) & (2) $\vec{F} = \vec{F}_e + \vec{F}_m$

$$= q\vec{E} + q(\vec{v} \times \vec{B})$$

$$= q(\vec{E} + (\vec{v} \times \vec{B}))$$

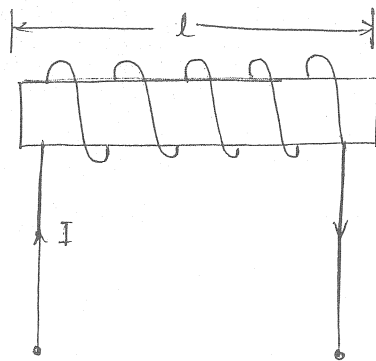
$$\boxed{\vec{F} = q(\vec{E} + (\vec{v} \times \vec{B}))}$$

Lorentz force equation

* It is used to find the total net force acting on a charge when Electric field, flux density is given whose charge is moving with specified velocity of q C of charge.

- (c) Calculate the inductance of solenoid of 200 turns wound tightly on cylindrical tube of length 60cm and diameter @ 6cm with air as media. derive the expression used

→ Inductance of solenoid



Consider a solenoid of N number turns as shown in above fig let current flowing through solenoid I amps let l be the length of solenoid and A be the cross sectional area Field intensity inside the solenoid

$$H = \frac{NI}{l} \text{ A/m} \quad \text{--- (1)}$$

The total flux linkage is given by

$$\text{Total flux linkage} = N\phi = N(B)(A)$$

$$\phi = N(\mu H)A$$

$$\text{Total flux linkage } \phi = \mu NHA$$

$$\text{we have } H = \frac{NI}{l} \text{ From (1)}$$

$$\therefore \text{Total flux linkage } \phi = \mu N \left(\frac{NI}{l} \right) A$$

$$\therefore \phi = \frac{\mu N^2 IA}{l}$$

$$\text{Inductance of solenoid } L = \frac{\text{Total flux linkage}}{\text{Total } C/n}$$

$$\therefore L = \frac{\mu N^2 I A}{l I}$$

$$\therefore \boxed{L = \frac{\mu N^2 A}{l}} \quad \text{or}$$

$$\boxed{L = \frac{\mu_0 \mu_r N^2 A}{l}}$$

Given : $N = 200$ turns

$$\mu = \mu_0 = 4\pi \times 10^{-7} \text{ Wb/A}\cdot\text{m}$$

$$d = 6\text{cm} = 6 \times 10^{-2} \text{ m} \quad \text{hence } r = \frac{d}{2} = 3 \times 10^{-2} \text{ m}$$

$$l = 60\text{cm} = 60 \times 10^{-2} \text{ m}$$

The inductance of solenoid is

$$L = \frac{\mu N^2 A}{l}$$

$$= \frac{4\pi \times 10^{-7} \times (200)^2 \times \pi (3 \times 10^{-2})^2}{60 \times 10^{-2}}$$

$$L = 2.3687 \times 10^{-4} \text{ H}$$

$$\boxed{L = 0.2368 \text{ mH}}$$

6(a) With neccessity relation ships, Explain Faraday's law's electromagnetic induction for both static and time varying feilds.

Faraday's 1st Law

Whenever a changing flux linking with the coil an. emf gets induced in the coil is known as Faraday's 1st law

Faraday's IInd Law

Magnitude of Induced emf is equal to the rate of change of Flux linkage

Faraday's IIIrd Law (Lenz's Law)

"The emf induced in the coil is opposite to own cause for it" the (-) minus sign represent the induced emf is in opposite direction of applied voltage.

$$e \propto -\frac{d\phi}{dt}$$

$$e = -N \frac{d\phi}{dt} \quad \text{where } N = \text{no. of conductors}$$

$$\oint \vec{E} \cdot d\vec{l} = -\frac{d}{dt} \int_S \vec{B} \cdot d\vec{s}$$

$$\boxed{\oint \vec{E} \cdot d\vec{l} = -\frac{\partial}{\partial t} \int_S \vec{B} \cdot d\vec{s}} \quad \text{① Integral form of Faradays Law}$$

From Stokes theorem we have

$$\oint \vec{H} \cdot d\vec{l} = \int_S (\nabla \times \vec{H}) \cdot d\vec{s}$$

$$\oint \vec{E} \cdot d\vec{l} = \int_S (\nabla \times \vec{E}) \cdot d\vec{s}$$

$$\int_S (\nabla \times \vec{E}) \cdot d\vec{s} = \frac{\partial}{\partial t} \int_S \vec{B} \cdot d\vec{s} \quad \because \text{From ①}$$

$$\boxed{\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}}$$

is called Faraday's Law in point form or
Maxwell's Ist equation.

(b) Starting From Faraday's Law of electromagnetic Induction

derive $\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$

→ From Faraday's law of induction "Induced emf is equal to the rate of change of magnetic flux linking with the closed circuit"

$$e = -N \frac{d\phi}{dt} \text{ volts} \text{ --- (1)}$$

where $N =$ number of turns

$e =$ induced emf

Let us assume single turn circuit for $N=1$

$$\therefore e = -\frac{d\phi}{dt} \text{ volts} \text{ --- (2)}$$

The minus sign in both eqⁿs (1) & (2) indicates the direction of induced emf in the coil is ~~very~~ opposite to very own cause for it is applied voltage which is Lenz's Law

Let us consider the Faraday's law of induced emf is scalar quantity measured in volts thus induced emf is given as

$$e = \oint \vec{E} \cdot d\vec{l} \text{ --- (3)}$$

The above equation is voltage about closed path

The magnetic flux passing through the surface area is

$$\phi = \int_S \vec{B} \cdot d\vec{s}$$

where $B =$ magnetic flux density

eqⁿ (2) can be rewritten as

$$e = -\frac{d}{dt} \int_S \vec{B} \cdot d\vec{s} \text{ --- (4)}$$

From eqⁿ (3) & (4) we get

$$e = \oint \vec{E} \cdot d\vec{l} = -\frac{d}{dt} \int_S \vec{B} \cdot d\vec{s} \text{ --- (5)}$$

These are two conditions for induced emf the closed ckt in which the induced emf is stationary varying with time From (5) \vec{B} may be changing coordinates as well as time

$$\oint \vec{E} \cdot d\vec{l} = - \int_s \frac{\partial \vec{B}}{\partial t} \cdot d\vec{s}$$

This is similar to the transformer action

$$\int_s (\nabla \times \vec{E}) \cdot d\vec{s} = - \int_s \frac{\partial \vec{B}}{\partial t} \cdot d\vec{s}$$

Assuming both the integrals are taken over closed identical surfaces

$$(\nabla \times \vec{E}) \cdot d\vec{s} = - \frac{\partial \vec{B}}{\partial t} \cdot d\vec{s}$$

Hence, Finally

$$\boxed{\nabla \times \vec{E} = - \frac{\partial \vec{B}}{\partial t}}$$

(e) Find the frequency at which conduction current density and displacement current density are equal in medium with

$$\sigma = 2 \times 10^4 \text{ S/m and } \epsilon_r = 81$$

→ we have the conduction current density equation as

$$\left| \frac{J_c}{J_D} \right| = \frac{\sigma}{\omega \epsilon} = \frac{2 \times 10^4}{2\pi f \times 81 \times \epsilon_0}$$

$$\therefore \boxed{f = 44.38 \text{ kHz}}$$

7 (a) What is uniform plane wave? Explain its propagation in free space with necessary equation

→ An electromagnetic wave travelling in a plane with uniform velocity is called as uniform plane wave.

$$\nabla \times \bar{E} = -\frac{\partial \bar{B}}{\partial t} = -\mu \frac{\partial \bar{H}}{\partial t} \quad \text{Maxwell's third equation}$$

$$\begin{aligned} \nabla \times \bar{H} &= \bar{J}_c + \bar{J}_c \\ &= \sigma \bar{E} + \epsilon \frac{\partial \bar{E}}{\partial t} \quad \text{--- (2)} \end{aligned}$$

Taking curl on both sides

$$\nabla \times (\nabla \times \bar{E}) = -\mu \nabla \times \frac{\partial \bar{H}}{\partial t}$$

$$\nabla \times (\nabla \times \bar{H}) = -\mu \frac{\partial}{\partial t} (\nabla \times \bar{H})$$

From vector Identity $\nabla(\nabla \cdot \bar{E}) - \nabla^2 \bar{E} = \nabla \times \nabla \times \bar{E}$

$$\nabla(\nabla \cdot \bar{E}) - \nabla^2 \bar{E} = \mu \frac{\partial}{\partial t} (\sigma \bar{E} + \epsilon \frac{\partial \bar{E}}{\partial t})$$

$$\nabla(\nabla \cdot \bar{E}) = \nabla^2 \bar{E} = -\mu \sigma \frac{\partial \bar{E}}{\partial t} - \mu \epsilon \frac{\partial^2 \bar{E}}{\partial t^2}$$

$$\nabla(\nabla \cdot \bar{E}) = \nabla^2 \bar{E} - \mu \sigma \frac{\partial \bar{E}}{\partial t} - \mu \epsilon \frac{\partial^2 \bar{E}}{\partial t^2}$$

$$\nabla^2 \bar{E} - \mu \sigma \frac{\partial \bar{E}}{\partial t} - \mu \epsilon \frac{\partial^2 \bar{E}}{\partial t^2} = \nabla \left(\frac{\rho_v}{\epsilon} \right)$$

For free space and charge free region

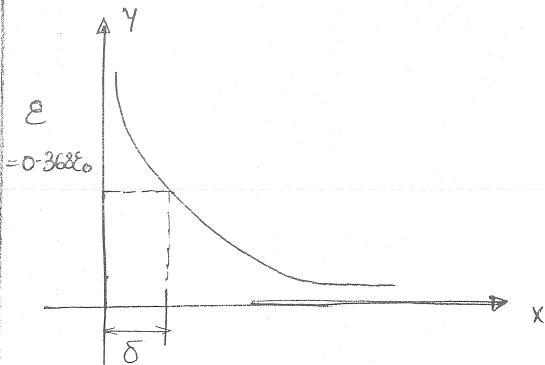
$$\boxed{\nabla^2 \bar{E} - \mu \epsilon \frac{\partial^2 \bar{E}}{\partial t^2} = 0}$$

Wave equation in free space

(b) Define Skin depth and depth of penetration.

→ Skin effect can be defined as the

Depth of Penetration



$$E = E_0 e^{-\alpha x}$$

$$\alpha \mu \sigma = \delta = \frac{1}{\alpha}$$

$$E = E_0 e^{-\alpha \frac{1}{\alpha}}$$

$$= E_0 e^{-1}$$

$$\therefore E = 0.3686 E_0$$

α for good conductor

$$\alpha = \sqrt{\frac{\mu \omega \epsilon}{2}}$$

$$\delta = \frac{1}{\alpha} = \frac{1}{\sqrt{\frac{\mu \omega \epsilon}{2}}}$$

$$\delta = \sqrt{\frac{2}{\omega \mu \epsilon}}$$

$$\delta = \sqrt{\frac{2}{2\pi f \mu \epsilon}}$$

$$\therefore \delta = \frac{1}{\sqrt{\pi f \mu \sigma}}$$

(c) For copper conductivity is 58 Ms/m find skin depth at a frequency of 10 MHz

→ Given $\sigma = 58 \text{ Ms/m}$

$$f = 10 \text{ MHz}$$

$$\mu_0 = 4\pi \times 10^{-7} \quad \mu_r = 1$$

Depth of Penetration or skin depth can be given as

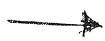
$$\delta = \frac{1}{\sqrt{\omega \mu \sigma}}$$

$$= \frac{1}{\sqrt{3.142 \times 10 \times 10^6 \times 4\pi \times 10^{-7} \times 58 \times 10^6}}$$

$$= 2.0895 \times 10^{-5}$$

$$\delta = 20.89 \mu\text{m}$$

8 (a) With necessary equation explain standing wave equation ratio



Consider a uniform plane wave travelling in lossless medium with no reflected wave present in medium. It is observed that relative amplitude of field same at every point. It is obvious that instantaneous field undergoes phase change of $\beta(z_2 - z_1)$ when the probe is moved from $z = z_1$ to $z = z_2$ along a wave thus we can conclude an unattenuated travelling wave shows equal voltage amplitude characteristics.

Now consider uniform plane wave travelling in lossless medium it gets reflected by back the perfect conductor. The result in generating standing wave equation when voltage probe is located at interference of $z = 0$ & at every integral no. of half wavelength from the interference in medium $\pm (z < 0)$. The o/p of the probe will be zero when positions of probe is interchanged the amplitude of field varies as $\sin |\beta z|$ as shown in fig.

- (b) Find whether the wet marshy soil characterised by $\rho = 10^3 \text{ s/m}$, $\epsilon_r = 15$ and $\mu = 1$ may be considered as conductor or dielectric or neither for the frequencies i) 60 Hz, ii) 1 MHz, iii) 100 MHz, (iv) 10 GHz

→ $\frac{\sigma}{\omega \epsilon} \gg 1$ For good conductor

$\frac{\sigma}{\omega \epsilon} \ll 1$ For good lossy dielectric conductor

$\frac{\sigma}{\omega \epsilon} = 0$ For perfect dielectric

(i) 60 Hz = f

$$\frac{\sigma}{\omega \epsilon} = \frac{10^{-2}}{2\pi \times 60 \times 15 \times \epsilon_0} = 20.98 \times 10^3 \gg 1$$

It is good conductor at 60 Hz

(ii) f = 1 MHz

$$\frac{\sigma}{\omega \epsilon} = \frac{10^{-2}}{2\pi \times 1\text{M} \times 15 \times \epsilon_0} = 11.98 \gg 1$$

Hence it is good conductor

(iii) f = 100 MHz

$$\frac{\sigma}{\omega \epsilon} = \frac{10^{-2}}{2\pi \times 100\text{M} \times 15 \times \epsilon_0} = 0.498 \ll 1$$

It is lossy dielectric conductor

(iv) f = 10 GHz

$$\frac{\sigma}{\omega \epsilon} = \frac{10^{-2}}{2\pi \times 10\text{G} \times 15 \times \epsilon_0} = 0.0011 \approx 0$$

For this frequency soil is perfect.

